

Feature-energy duality of topological boundary states in multilayer quantum spin Hall insulator



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(前沿量子科技研究中心)



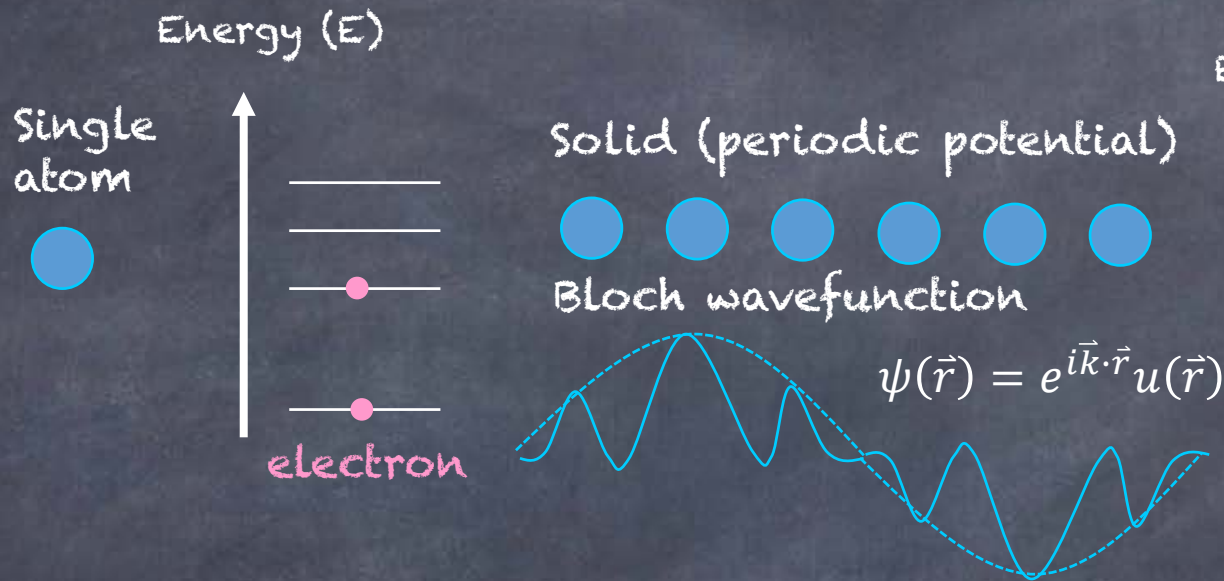
Physics Division, National Center for Theoretical Sciences
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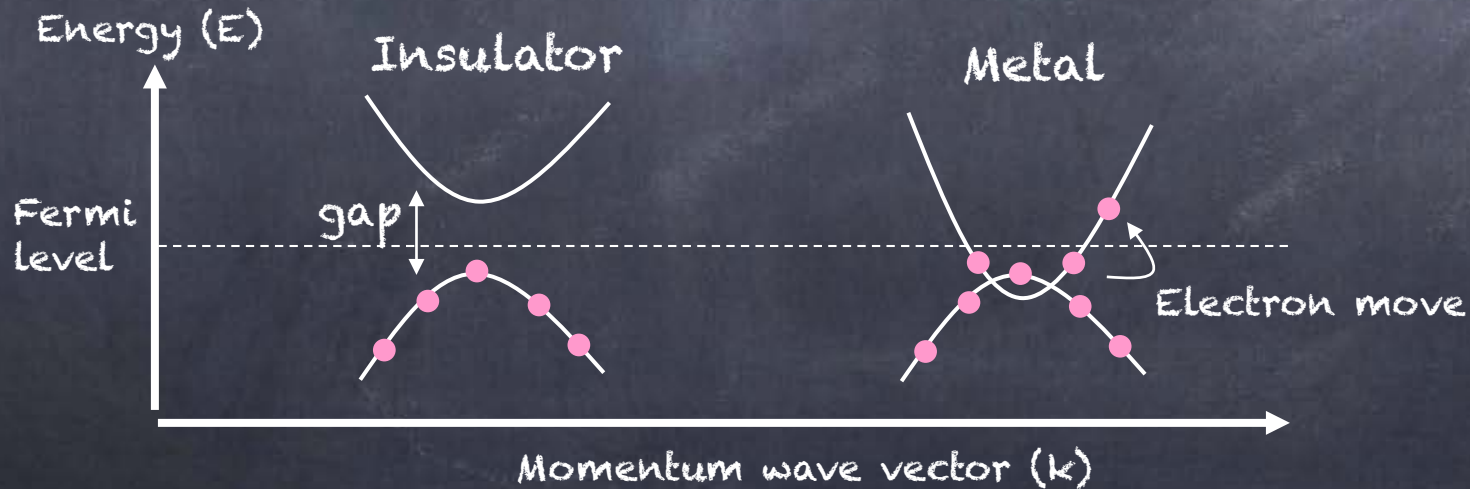
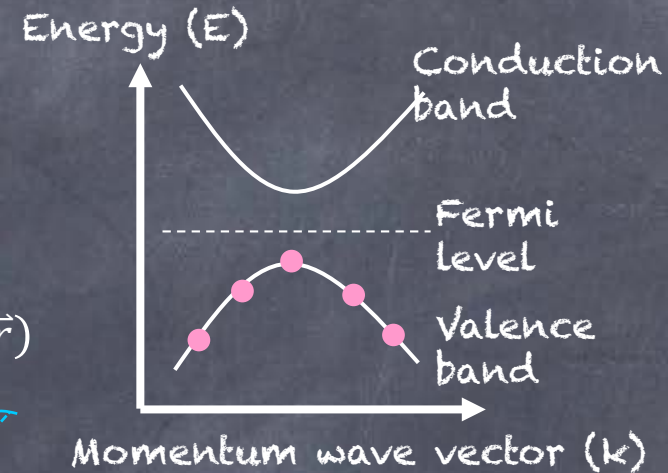
Tay-Rong Chang (張泰榕)

2024/Aug./28

Band theory



Band structure



Band theory (topology)

Math \Rightarrow real space

Phys \Rightarrow momentum space

Gauss-Bonnet Theorem:

$$\int_S K_{Gauss} ds = 4\pi(1 - g)$$

↑
Gauss curvature
↑
genus

TKNN theorem:

$$\frac{1}{2\pi} \sum_{m=occ} \int_{BZ} d^2k \underbrace{\vec{\nabla} \times i \langle u_m | \vec{\nabla} | u_m \rangle}_{\text{Berry curvature}} = \underbrace{n}_{\text{Topological invariant number}}$$

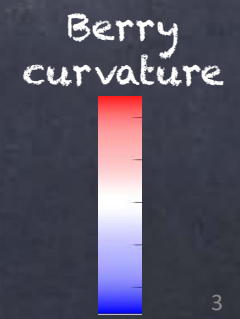
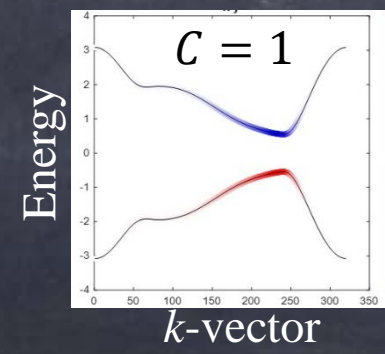
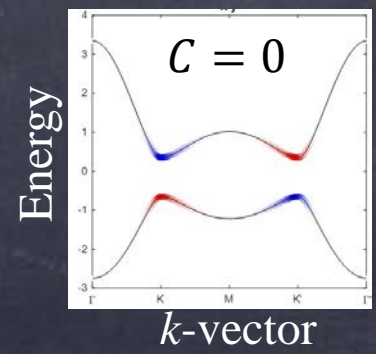
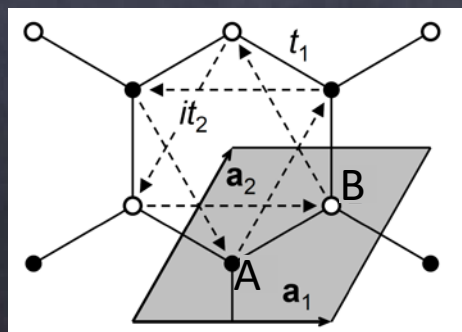
↙
Bloch wavefun

Haldane model (2D honeycomb lattice)

Ref: PRL61, 2015 (1988)

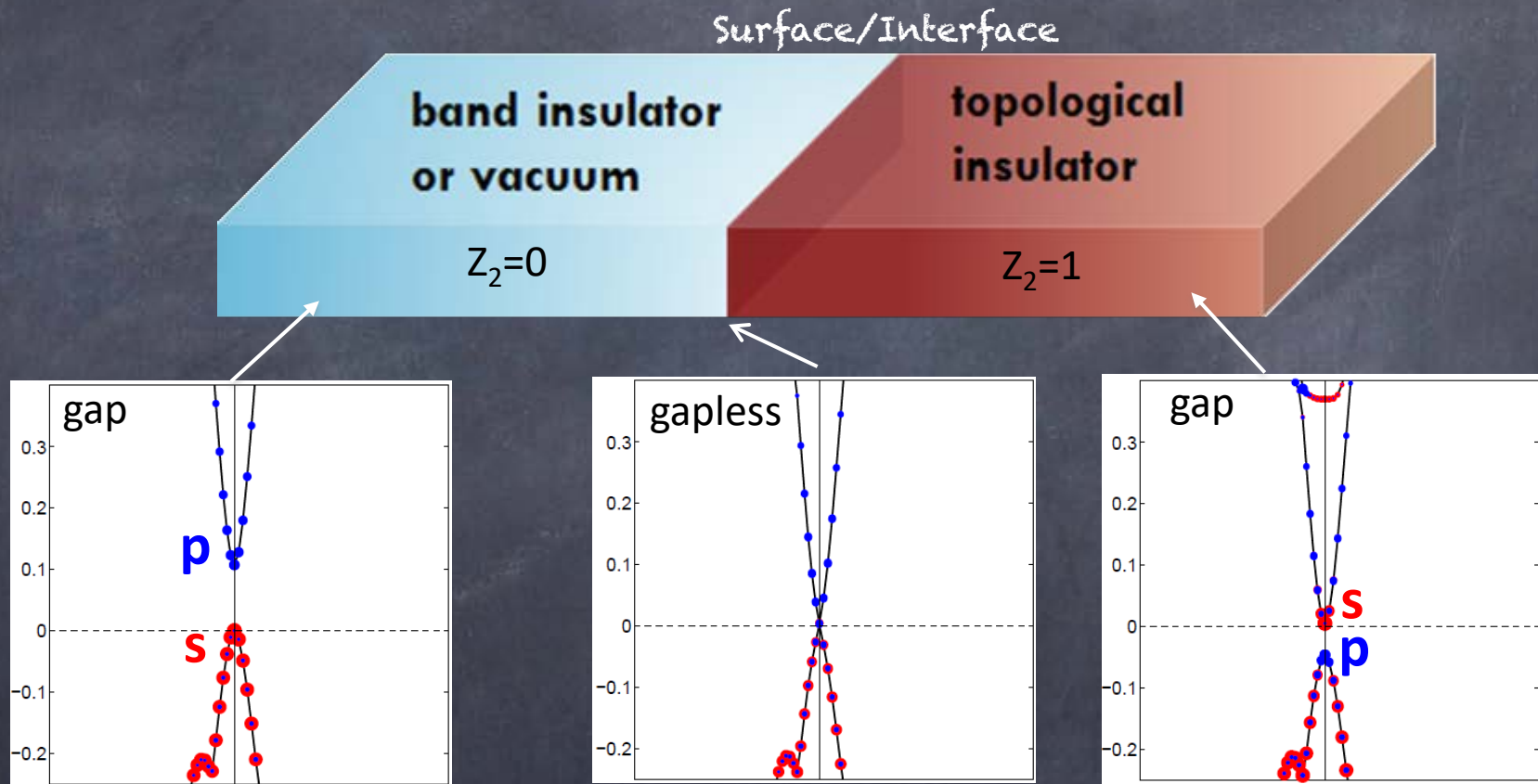
$$H = \sum_i (-1)^{\tau_i} M c_i^\dagger c_i + t_1 \sum_{\langle ij \rangle} (c_i^\dagger c_j + h.c.) + t_2 \sum_{\langle\langle ij \rangle\rangle} (e^{i\phi_{ij}} c_i^\dagger c_j + h.c.)$$

$M = 0.7, t_1 = -1, t_2 = 0$ $M = 0.7, t_1 = -1, t_2 = -0.24$



Band theory (topology)

Bulk-boundary correspondence



The gapless boundary state is the hallmark of topological phase

Ref: Rev. Mod. Phys. **82**, 3045 (2010)
Ref: Rev. Mod. Phys. **83**, 1057 (2011)

Topological non-trivial state

Bulk property

Topological invariant number

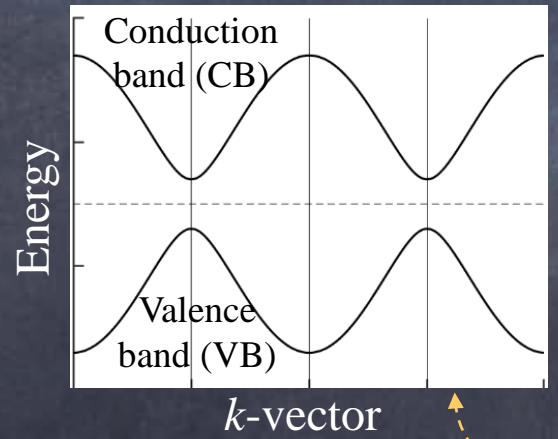
Chern number, \mathbb{Z}_2 index, \mathbb{Z}_4 index, winding number...etc

Bulk-boundary correspondence

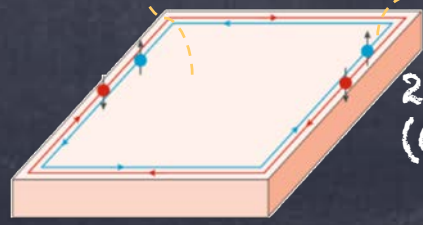
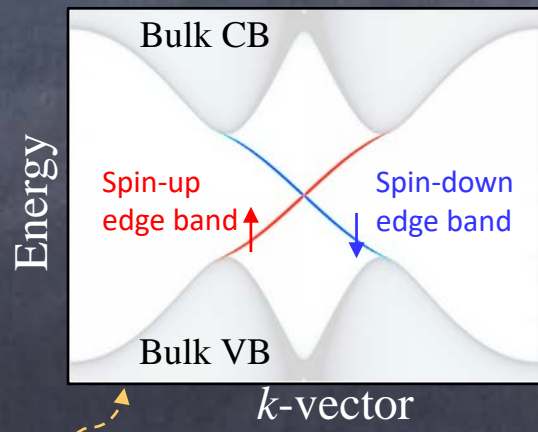
Edge property

Gapless topological boundary state

Topological invariant number: \mathbb{Z}_2 index ($\mathbb{Z}_2 = 1$)



Gapless edge band

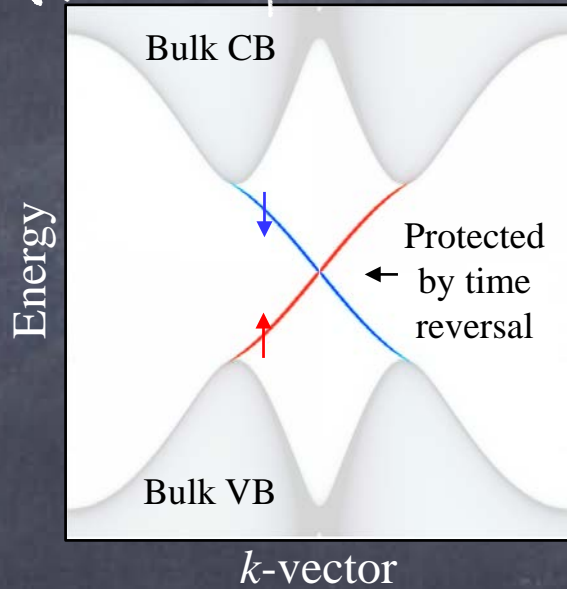


2D \mathbb{Z}_2 topological insulator (Quantum spin Hall insulator)

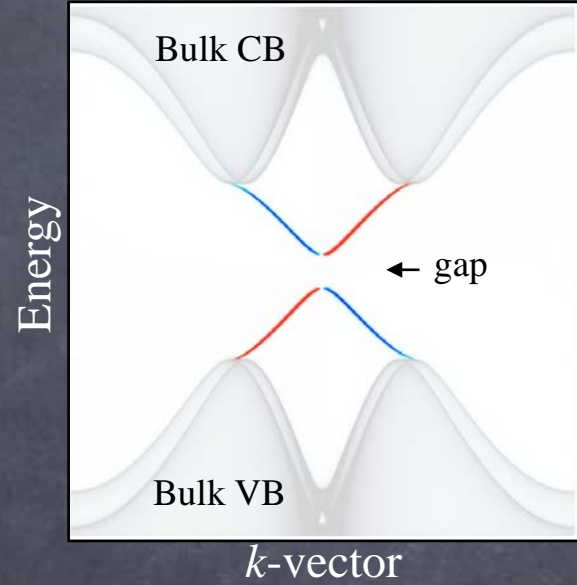
Topological non-trivial state

If we consider

2D \mathbb{Z}_2 topological insulator
(Quantum spin Hall insulator)



Symmetry breaking perturbation



Is this system topologically trivial?

- Band structure (symmetry): trivial
- Feature Spectrum Topology: non-trivial

Outline

- Topological nontrivial state (with) symmetry protection
2D \mathbb{Z}_2 topological insulator (Quantum spin Hall insulator)

The central question of this talk: If the boundary state no longer connects the valence band and the conduction band, is this system topologically trivial?

- Feature Spectrum Topology

Kane-Mele model

Feature-energy duality

- Material realization

Multilayer Bi_4Br_4

Kane-Mele model w/o Rashba SOC

Ref: PRL95, 226801 (2005) · PRL95, 146802 (2005)

Kane-Mele model (2D honeycomb lattice)

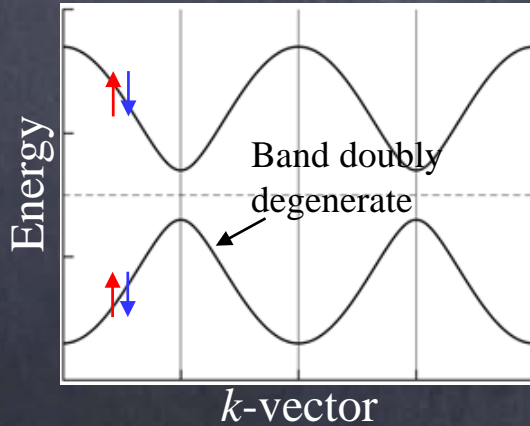
Matrix form

$$H_0 = t \sum_{\langle i,j \rangle \alpha} c_{i\alpha}^\dagger c_{j\alpha} + \underbrace{\frac{i\lambda_I}{3\sqrt{3}} \sum_{\langle\langle i,j \rangle\rangle \alpha\beta} v_{ij} c_{i\alpha}^\dagger (s_z)_{\alpha\beta} c_{j\beta}}_{\text{Intrinsic SOC}}$$

$$H_0 = \begin{pmatrix} \begin{matrix} |\uparrow\rangle \\ H_{QAH}^+ \\ 0 \end{matrix} & \begin{matrix} |\downarrow\rangle \\ 0 \\ H_{QAH}^- \end{matrix} \end{pmatrix} \begin{matrix} |\uparrow\rangle \\ |\downarrow\rangle \end{matrix} \text{Basis}$$

H_{QAH} Haldane Hamiltonian (2x2 matrix)

Band structure

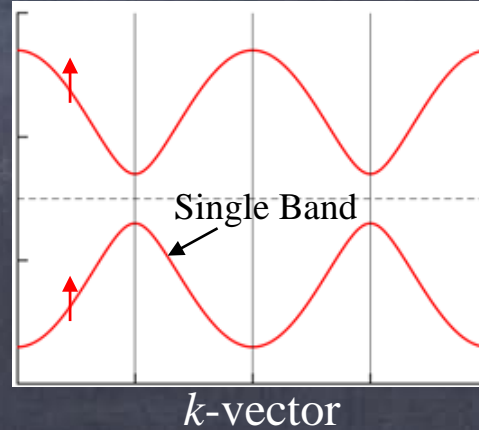


\hat{S}_z -conserved

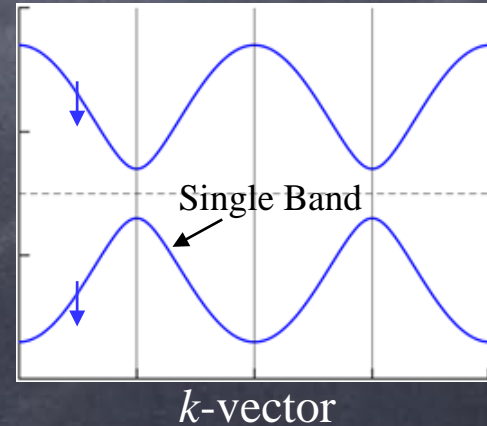


$|\uparrow\rangle_{VB}$ and $|\downarrow\rangle_{VB}$ separated

Spin-up (H_{QAH}^+)



Spin-down (H_{QAH}^-)

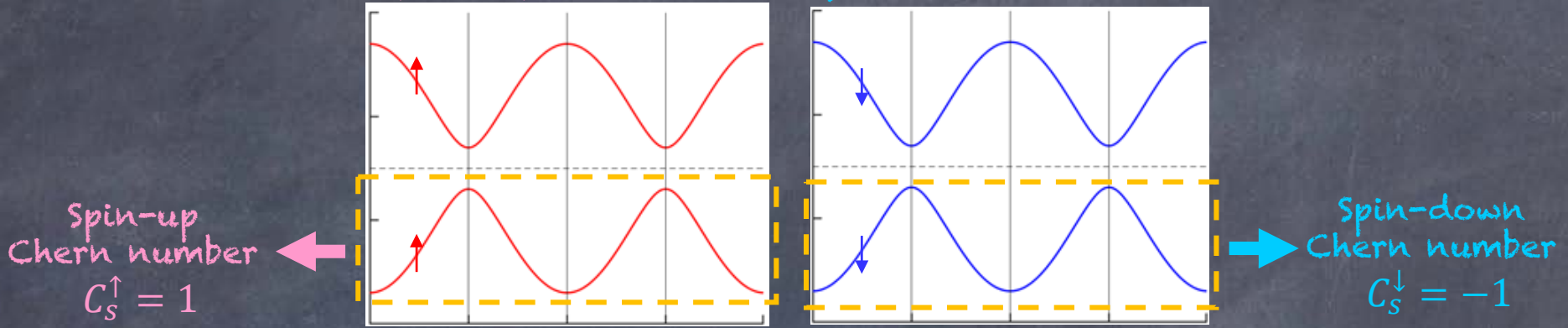


Kane-Mele model w/o Rashba SOC

Chern number: $C = \frac{1}{2\pi} \sum_{m=occ} \int dk^2 \vec{\nabla} \times i \langle u_m | \vec{\nabla} | u_m \rangle$ Bloch wavefun of VB

Spin-up (H_{QAH}^+)

Spin-down (H_{QAH}^-)



Spin Chern number (C_s)

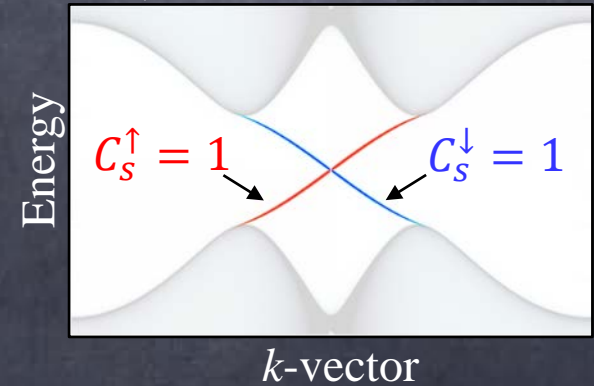
$$C_s = \frac{(C_s^\uparrow - C_s^\downarrow)}{2} = 1$$

$$\mathbb{Z}_2 = \text{mod}(C_s, 2) = 1$$

Bulk-boundary correspondence



Gapless edge bands



Kane-Mele model is a quantum spin Hall insulator

Kane-Mele model (\mathcal{T} -broken)

Kane-Mele model (2D honeycomb lattice)

$$H_M = H_0 - i \frac{2}{3} \lambda_{so} \sum_{\langle\langle i,j \rangle\rangle \alpha\beta} \mu_i c_{i\alpha}^\dagger (\mathbf{s} \times \hat{\mathbf{d}}_{ij})_{\alpha\beta}^z c_{j\beta} + M \sum_{i\alpha} c_{i\alpha}^\dagger s_z c_{i\alpha} - l \sum_{i\alpha} \mu_i c_{i\alpha}^\dagger c_{i\alpha}$$

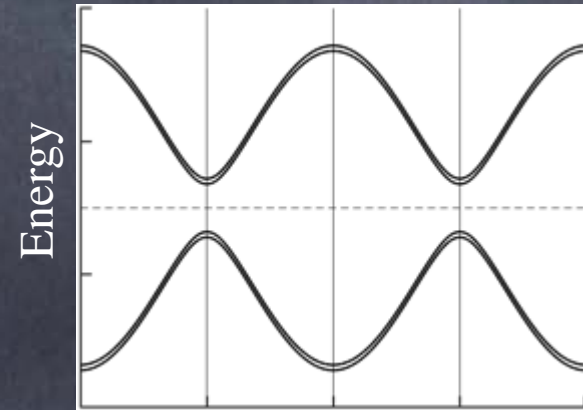
Kane-Mele (\hat{S}_z conserved) Rashba SOC: off-diagonal terms (break \hat{S}_z symmetry) Zeeman exchange field (break \mathcal{T})

Matrix form

$$H_M = \begin{pmatrix} H_{QAH}^+ + M & -H_R \\ H_R & H_{QAH}^- - M \end{pmatrix} \begin{matrix} |\uparrow\rangle \\ |\downarrow\rangle \end{matrix} \text{Basis}$$

H_{QAH} Haldane Hamiltonian Rashba SOC Zeeman

Band structure



k -vector
 $|\uparrow\rangle_{VB}$ and $|\downarrow\rangle_{VB}$ mixed

Kane-Mele model (\mathcal{T} -broken)

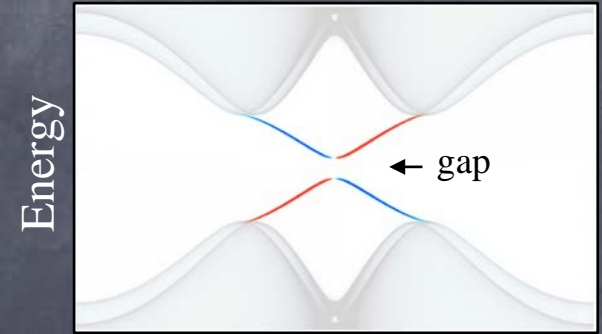
Matrix form

$$H_M = \begin{pmatrix} H_{QA}^+ + M & -H_R \\ H_R & H_{QA}^- - M \end{pmatrix} \begin{matrix} |\uparrow\rangle \\ |\downarrow\rangle \end{matrix}$$

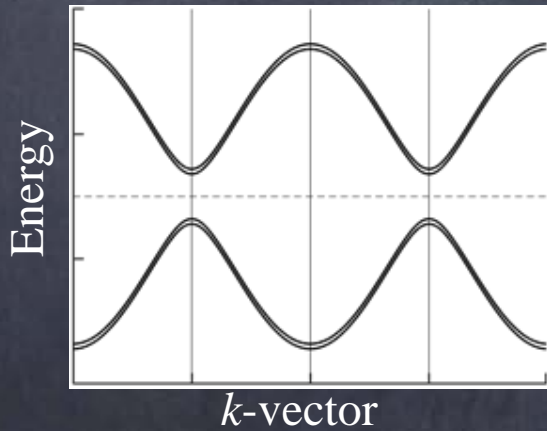
Basis

H_{QA}^+ Haldane Hamiltonian H_{QA}^- Haldane Hamiltonian
 H_R Rashba SOC M Zeeman

Edge bands

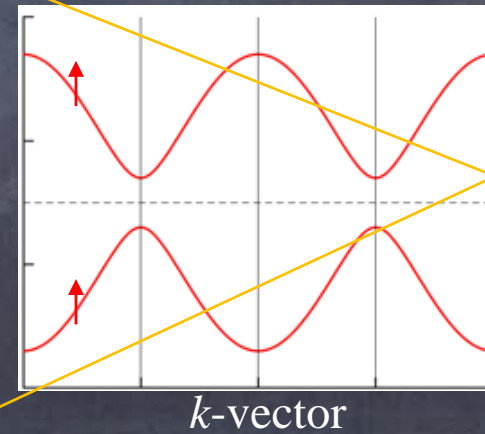


Band structure

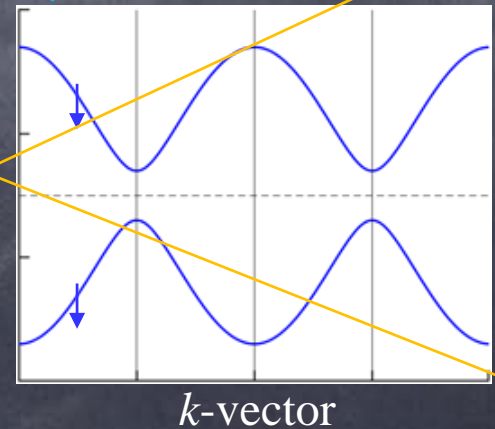


decompose

Spin-up (H_{QA}^+)



Spin-down (H_{QA}^-)



How to deal with topological phase in this model?

Feature Spectrum Topology

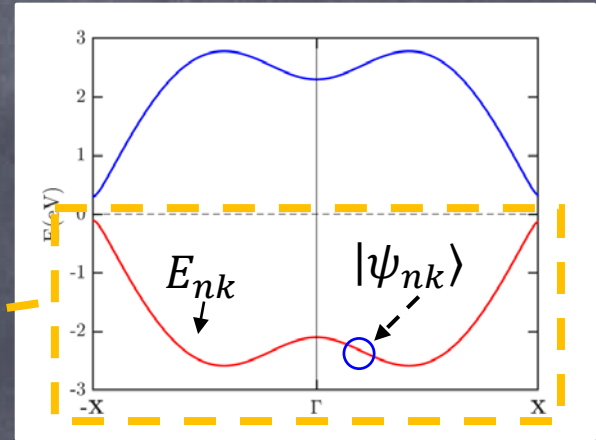
Ref: E. Prodan PRB 80, 125327 (2009), H. Lin arXiv:2310.14832v1 (2023)

Band structure:

$$\hat{H}|\psi_{nk}\rangle = E_{nk}|\psi_{nk}\rangle$$

Hamiltonian
eigenvalue
eigenstate

Band structure



$$C = \frac{1}{2\pi} \sum_{m=occ} \int dk^2 \Omega_m$$

The nature of band topology mainly comes from valence electron

Feature Spectrum Topology:

$$\hat{F}_O|\tilde{\psi}_{mk}\rangle = O_{mk}|\tilde{\psi}_{mk}\rangle$$

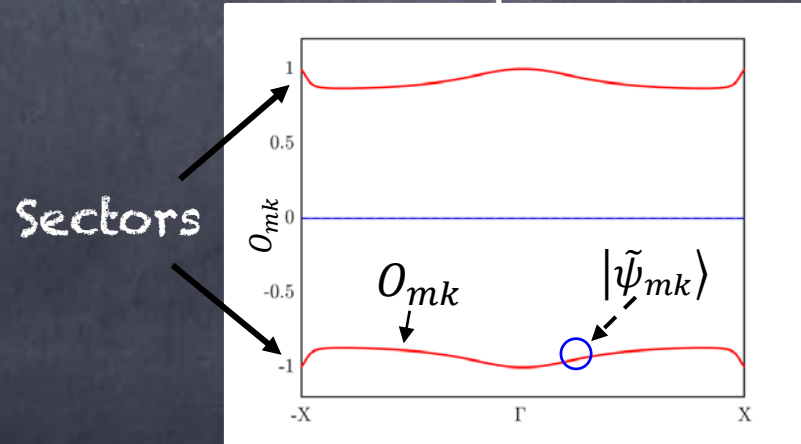
Feature operator
Feature eigenvalue
Feature eigenstate

$$\hat{F}_O = P\hat{O}P$$

Quantum operator
Projector (project to the occupied state)

$$|\tilde{\psi}_{mk}\rangle = \sum_{i=occ} a_i |\psi_i\rangle_{VB}$$

Feature spectrum



Topological invariant number: Feature Chern number

Feature Spectrum Topology

Kane-Mele model

\hat{H} (4x4 matrix)

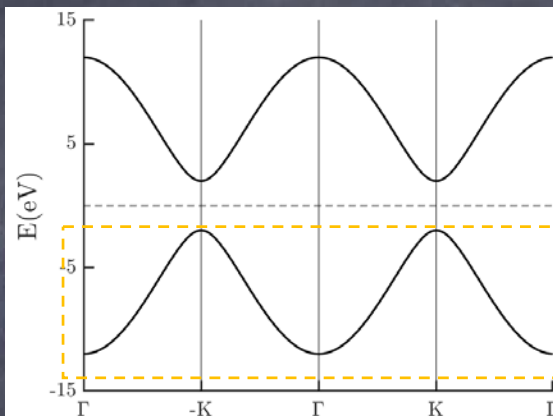
Projected \hat{S}_z onto the VB of the KM model

Feature operator

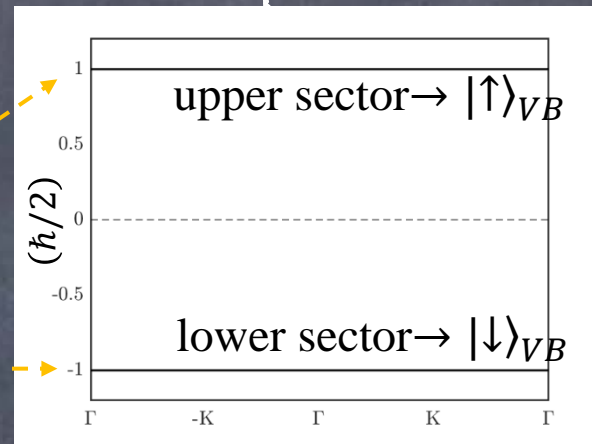
$\hat{F}_{S_z} = P\hat{S}_zP$ (2x2 matrix)



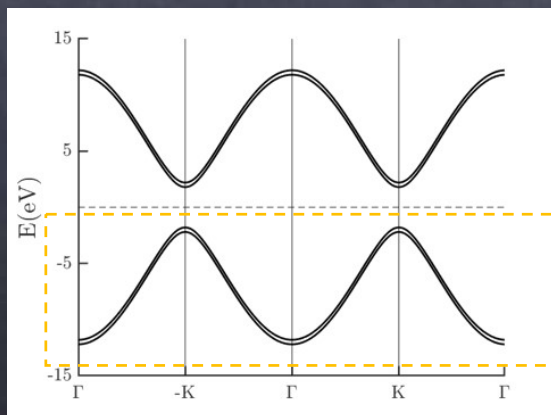
Band structure (w/o Rashba SOC)



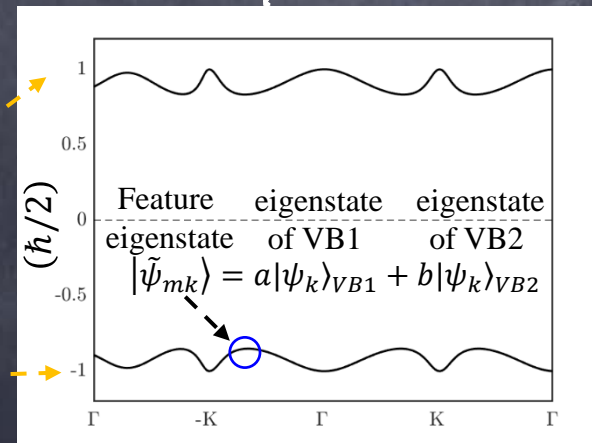
Feature spectrum



Band structure (w/ Rashba SOC, w/o T)



Feature spectrum



Feature Spectrum Topology

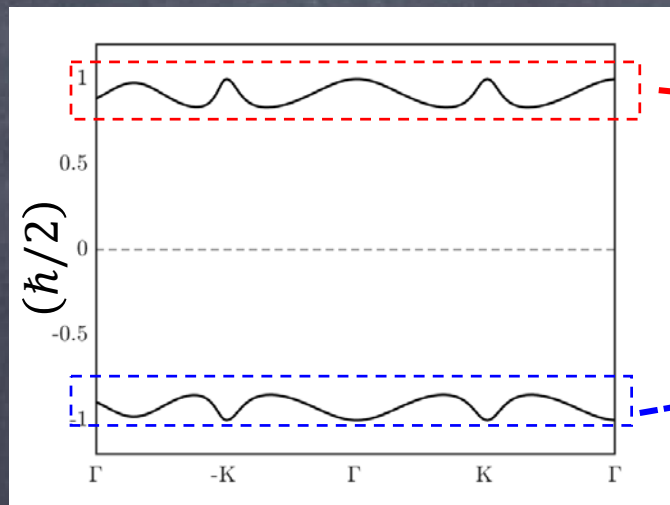
PRB109, 155143 (2024)

Feature operator

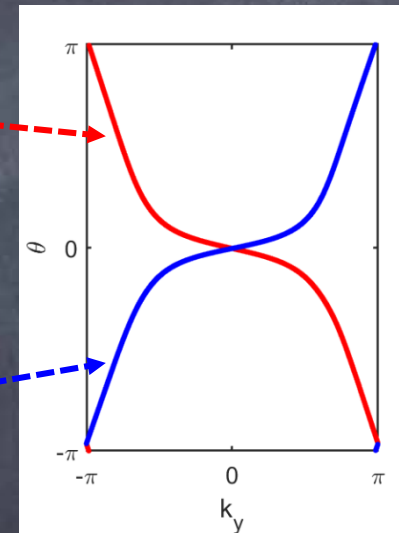
$$\hat{F}_{S_z} = P\hat{S}_zP$$



Feature spectrum



Wilson Loop



$$C_s^+ = 1$$

$$C_s^- = -1$$

Spin Chern number (C_s)
from feature eigenstates
(feature Chern number) $= C_s = \frac{(C_s^+ - C_s^-)}{2} = 1$

Kane-Mele model (with Rashba SOC and break \mathcal{T}) is still a QSHI

Feature Spectrum Topology

Bulk property

Topological invariant number

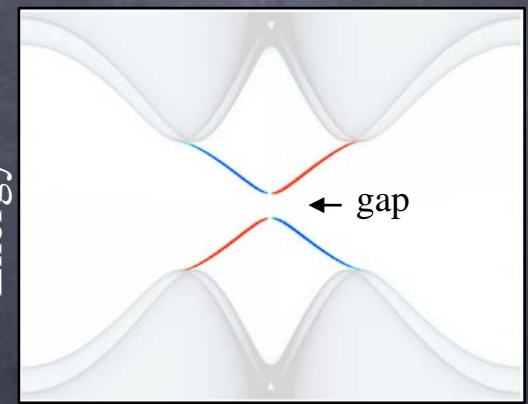
Bulk-boundary correspondence

Edge property

Gapless topological boundary state

Spin Chern number
(feature Chern number)

$$C_s = \frac{(C_s^+ - C_s^-)}{2} = 1$$

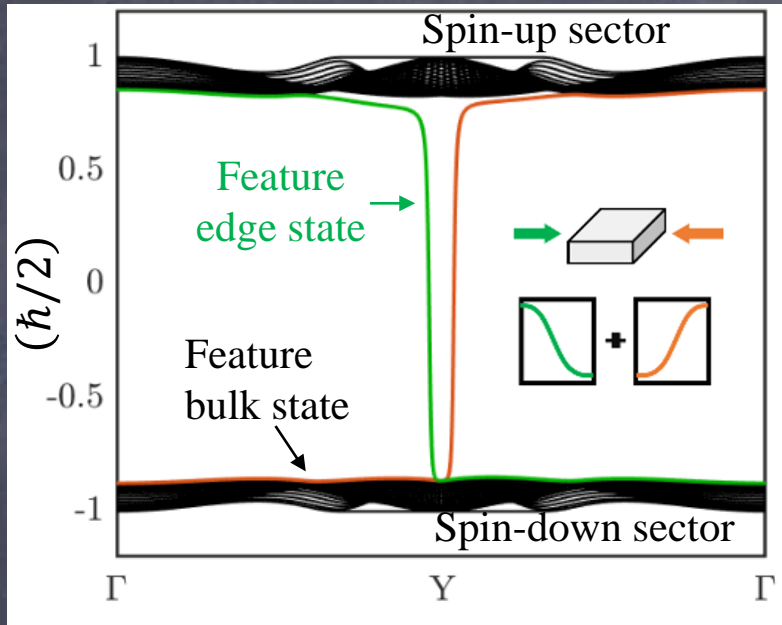


k-vector

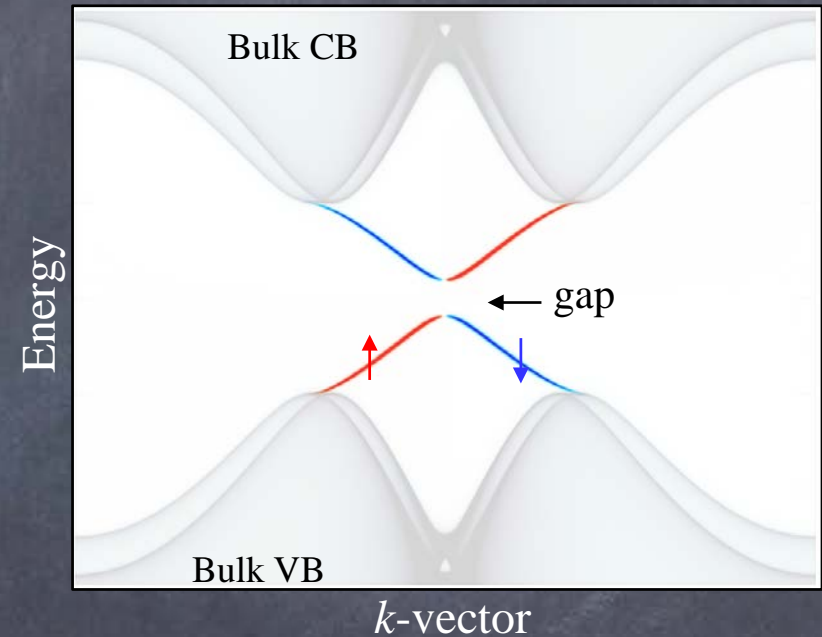
Edge Feature Spectrum

PRB109, 155143 (2024)

Edge feature bands



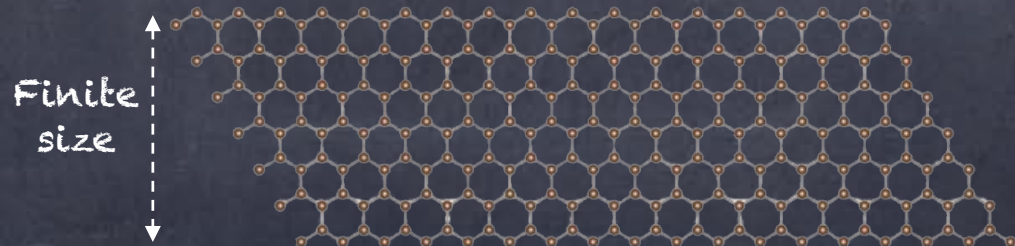
Edge bands (w/o T)



Spin Chern number
(feature Chern number)

$$C_s = \frac{(C_s^+ - C_s^-)}{2} = 1$$

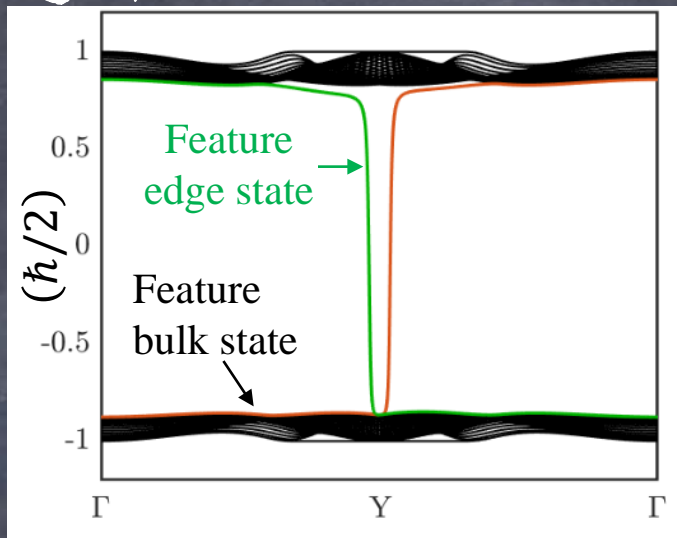
1D ribbon structure



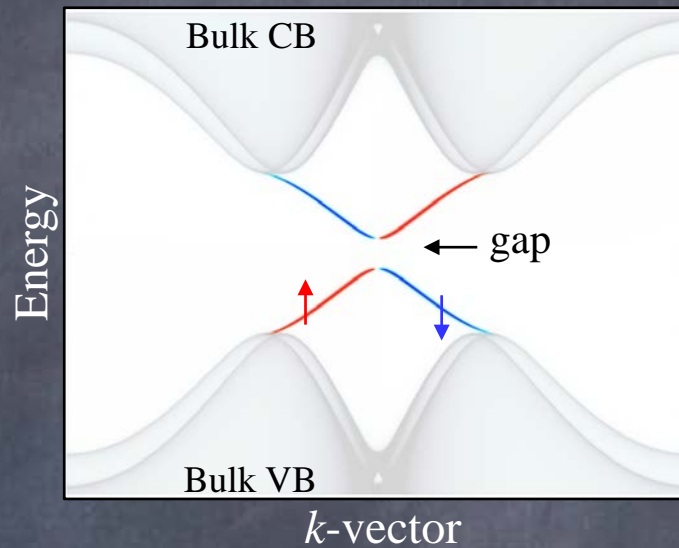
Edge Feature Spectrum

PRB109, 155143 (2024)

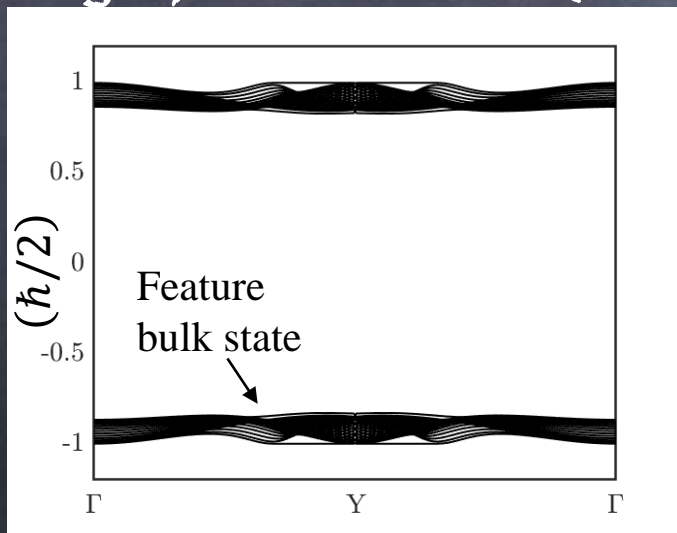
Edge feature bands (w/o \mathcal{T})



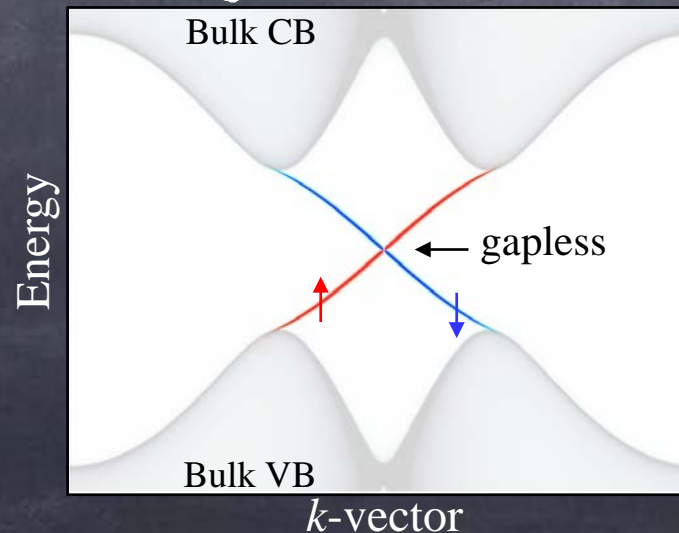
Edge bands (w/o \mathcal{T})



Edge feature bands (w/ \mathcal{T})



Edge bands (w/ \mathcal{T})



Feature-energy duality

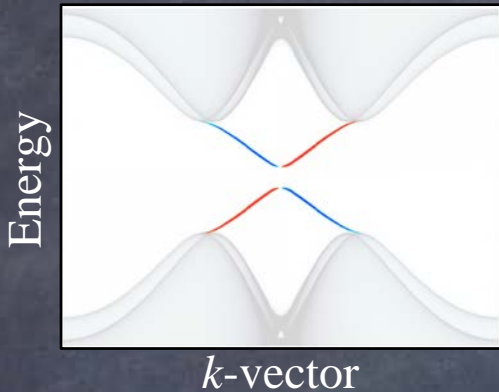
PRB109, 155143 (2024)

Band
(w/o \mathcal{T})

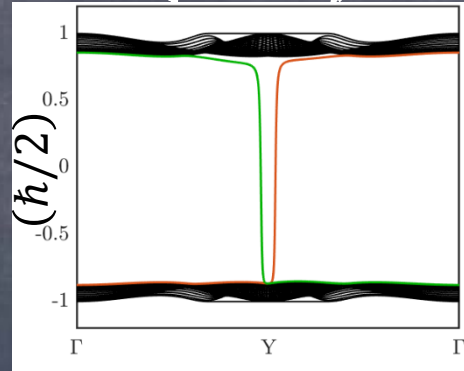
of gapless
edge states

Feature
(w/o \mathcal{T})

of gapless
edge states



0



1

Spin Chern
number

$$C_S = 1$$

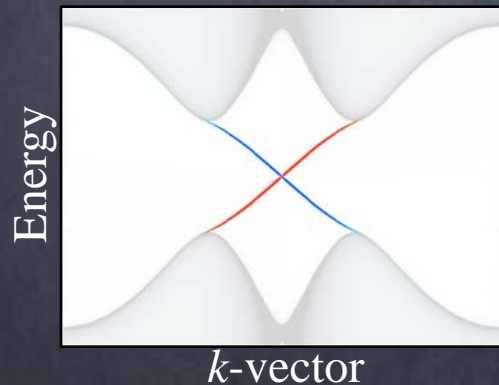
or

(w/ \mathcal{T})

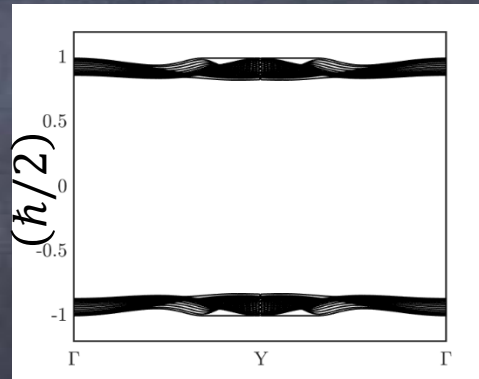
+

(w/ \mathcal{T})

=



1



0

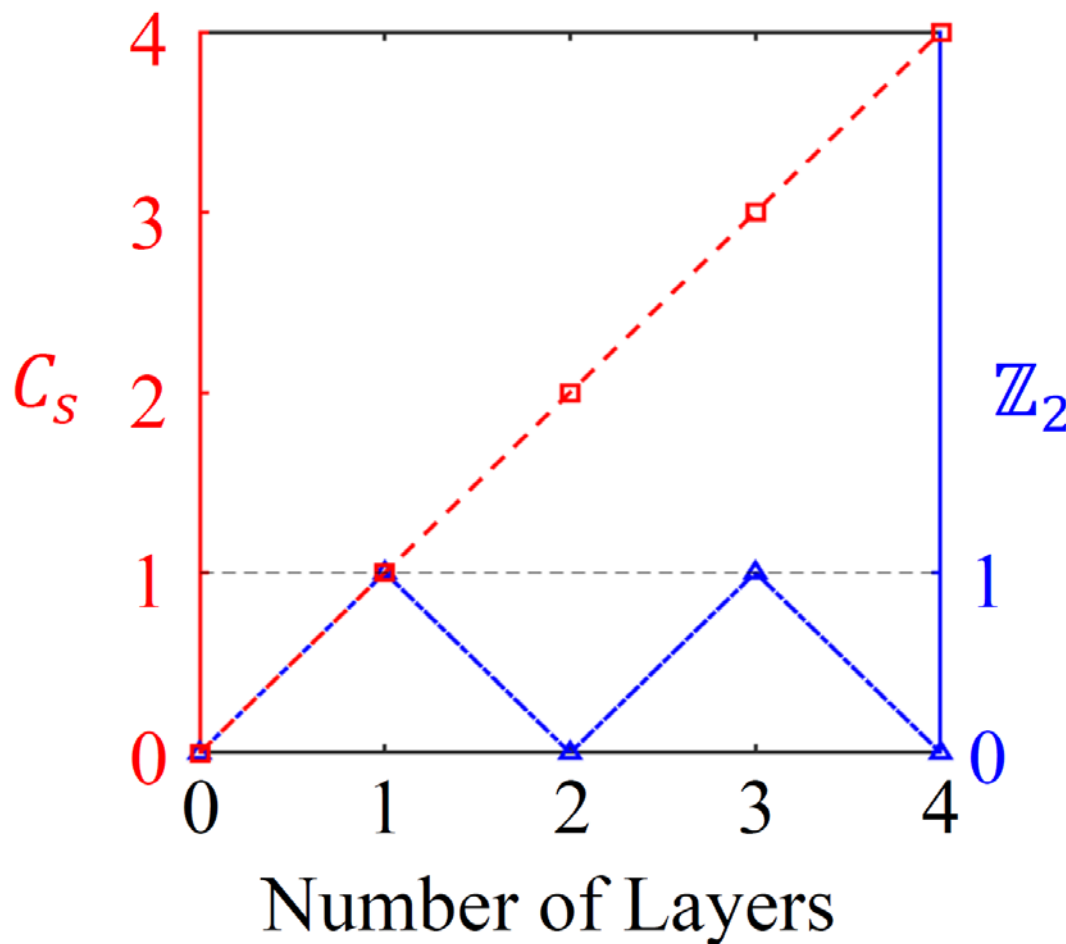
$$C_S = 1$$

Feature-energy duality: Total number of gapless edge states on the band structure, combined with the non-trivial edge states in the feature spectrum, equals the spin Chern number of QSHI.

Multilayer Kane-Mele model

PRB109, 155143 (2024)

Multilayer Kane-Mele model



Feature eigenstates:

Spin Chern number C_s exhibits linear increases with the number of layers.

Block eigenstates:

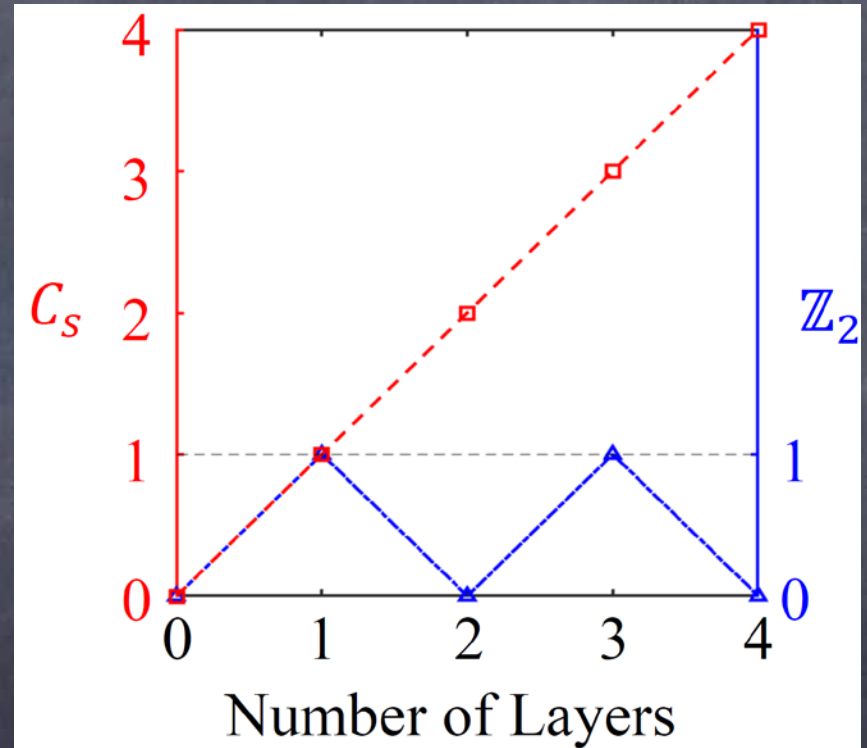
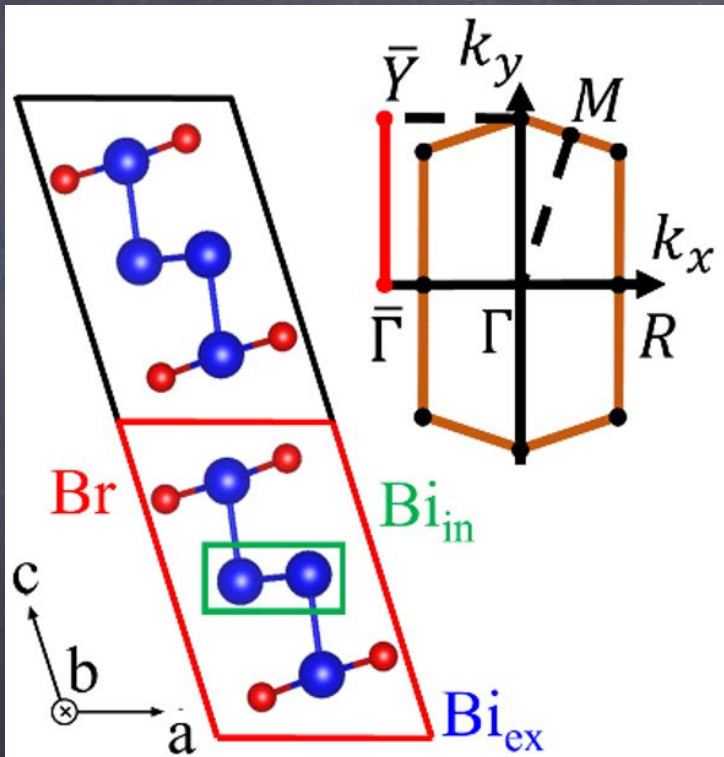
\mathbb{Z}_2 index oscillates between 1 and 0 for odd and even number of layers, respectively.

Material realization Bi_4Br_4

PRB109, 155143 (2024)

Quasi-1D structure

Same as multilayer Kane-Mele model



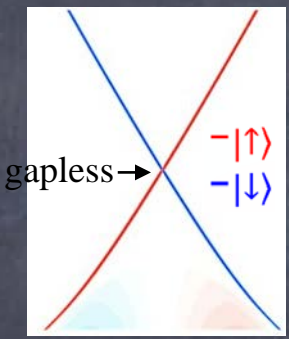
Material realization Bi_4Br_4

PRB109, 155143 (2024)

Monolayer Bi_4Br_4

$C_s = 1, \mathbb{Z}_2 = 1$

Bands



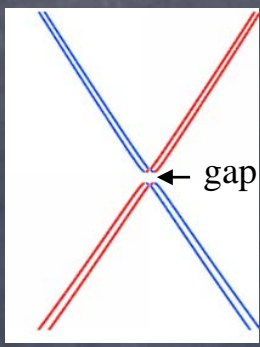
of gapless edge states

1

Bilayer Bi_4Br_4

$C_s = 2, \mathbb{Z}_2 = 0$

of gapless edge states

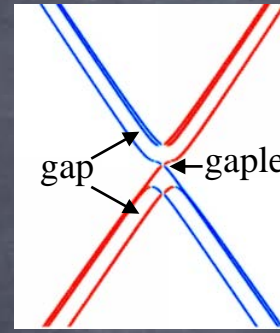


0

Trilayer Bi_4Br_4

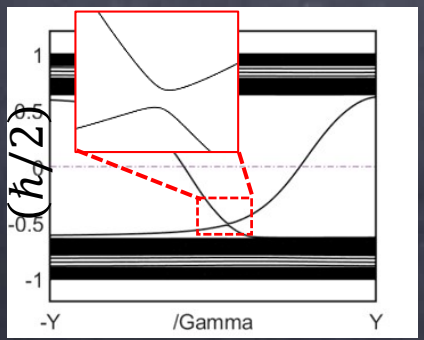
$C_s = 3, \mathbb{Z}_2 = 1$

of gapless edge states

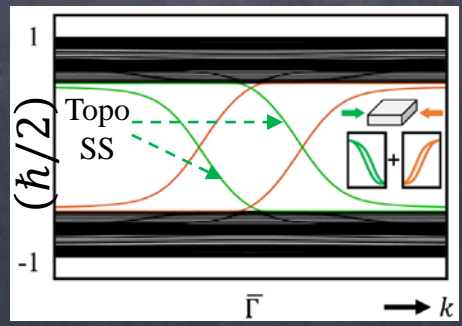


1

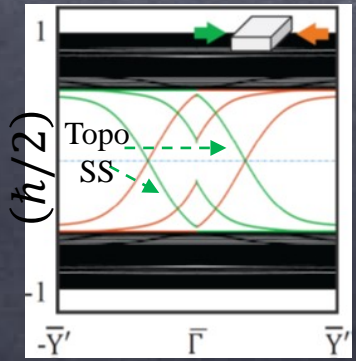
Feature



0



2



2

Conclusion

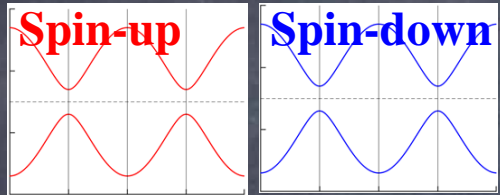
Kane-Mele model
(\hat{S}_z conserved)

Rashba SOC
w/o T

Kane-Mele model
(w/ Rashba SOC)
(w/o T)

$$\hat{F}_{S_z} = P \hat{S}_z P$$

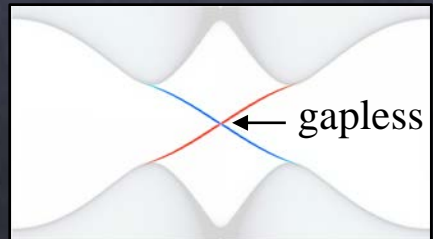
Feature operator



Spin Chern number
(Bloch eigenstate)

$$C_s = \frac{(C_s^\uparrow - C_s^\downarrow)}{2} = 1$$

Gapless boundary states
in the band structure

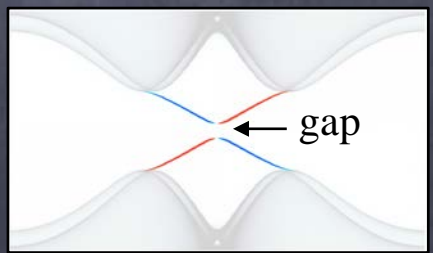


S_z is no longer a
good quantum number

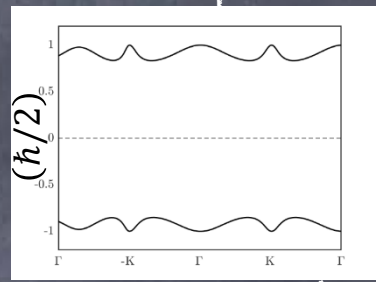
~~Spin Chern number
(Bloch eigenstate)~~

~~$$C_s = \frac{(C_s^\uparrow - C_s^\downarrow)}{2} = 1$$~~

Gap boundary states in
the band structure



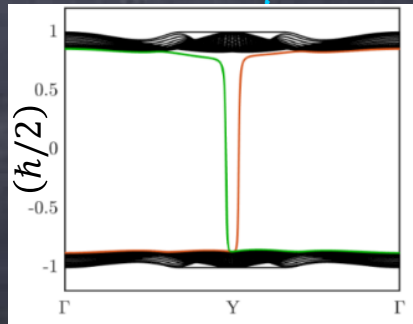
Feature spectrum



Spin Chern number
(feature eigenstate)

$$C_s = \frac{(C_s^\uparrow - C_s^\downarrow)}{2} = 1$$

Gapless boundary states
in the feature spectrum



Feature-energy duality

Acknowledgements

Feature-energy duality of topological boundary states in a multilayer quantum spin Hall insulator

Phys. Rev. B 109, 155143 (2024)
arXiv:2312.11794 (2023)

My master student

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Collaborators

Hsin Lin (Sinica)



Yi-Chun Hung
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Feature Spectrum Topology
arXiv:2310.14832 (2023)

Thank you!

To see a world in a grain of sand ... —William Blake

