Air-Bearing Effects on Actuated Thermal Pole-Tip Protrusion for Hard Disk Drives

Jia-Yang Juang
David B. Bogy

Computer Mechanics Laboratory, Department of Mechanical Engineering, University of California at Berkeley, Berkeley, CA 94720

1 Introduction

With the increase of areal density in hard disk drives the physical spacing (or flying height (FH)) between the read/write element and the surface of the disk has been continuously decreased. A spacing of ~2.5 nm is said to be required for a density of 1 Tbit/in². At such a low FH, static losses of the FH due to manufacturing tolerance, ambient pressure changes, and temperature variations can cause head-disk contact and result in data loss. Furthermore, slider disk contacts must be avoided during load/unload processes and operational shocks. The dynamic instability caused by FH modulations (FHMs) and nanoscale adhesion forces, such as electrostatic and intermolecular forces, should be minimized. Those challenges make a conventional air-bearing surface (ABS) slider an unlikely choice for a Tbit/in². One potential solution is a FH adjustment or controlled slider that is capable of adjusting its gap FH. Because of their quick response and low power consumption, piezoelectric materials have been proposed as active elements for adjusting the FH [1–8]. However, the requirements of the piezoelectric materials and the necessary modification of the slider design pose challenges in integration of the fabrication process and increase the manufacturing cost.

The read/write elements of a magnetic head slider consist of thin layers of different materials, including write poles, sensor, shields, undercoat, overcoat, coil, insulation layer, and the substrate. Because of the mismatch of the coefficients of thermal expansion of the various materials, the pole-tip of the write pole protrudes below the ABS plane when the ambient temperature varies and/or when an internal heat source is generated by Joule heating of the write current. These temperature-induced pole-tip protrusions (T-PTP) and write-induced pole-tip protrusions (W-PTP) adversely reduce the FH by several nanometers and increase the risk of head-disk contact. Several groups have numerically and/or experimentally investigated the effects of T-PTP and W-PTP [9–17].

Based on this concept, Meyer et al. [18] deposited a resistance heating element (heater) near the read/write elements, and the gap FH was reduced by applying a current through the heater to deliberately induce the pole-tip protrusion. Similarly, Kurita et al. developed an active head slider with a nanothermal actuator [19]. They used a finite element method to calculate the temperature distribution and thermal protrusion of their slider. They found that the additional air pressure increase caused by the protrusion lifted the slider upward and the amount of FH reduction was 30% less than the protrusion. In their study, the distribution of the heat transfer coefficient on the ABS was assumed to be constant. However, the effect of heat conducted from the slider to the disk through the ABS is a strong function of both FH and air pressure distributions, and hence, these two factors have to be considered in the model. Juang et al. [20] studied the actuation performance of an ABS slider with consideration of the effect of FH and pressure distributions. They found that even though the protruded area was relatively small, there was still considerable air bearing coupling with the resulting actuation efficiency (defined as the ratio of FH reduction to protrusion) of only 63%, which suggested that ABS played a key role in the actuation performance. Therefore, it is highly desirable to have a better understanding of the effects of the ABS on the thermal actuation and to provide design guidelines for improving the performance.

In this paper, we study the effects of the ABS on thermal actuation by numerical simulation. A three-dimensional (3D) thermal-structural coupled field finite element model is created with detailed structures of read/write and heating elements. The cooling effect of the air bearing is included in the model as ther-
2 Numerical Modeling and Analysis

The temperature distribution and thermal deformation of a slider body with thermal actuation are determined by the configuration, dimensions, and material properties of the read/write and heating elements and boundary conditions. The heat transfer boundary conditions depend on the flying state of the slider and, hence, the air bearing surfaces. In this study, we kept the structure of the transducers the same and investigated the effect of ABS on actuation performance. We created a three-dimensional finite element model of an entire slider (length=1.25 mm, width=1.00 mm, thickness=0.30 mm) with a detailed read/write transducer structure as shown in Fig. 1. The slider has a single-layer five-turn copper coil, a yoke width of 12 μm, and a write track width of 1 μm. The top and bottom poles are 1 μm thick. The top and bottom magnetic shields are 2 μm thick. The heating element has a thickness of 250 nm and is located between the coil and bottom pole. The photoresist layer, undercoat insulation layer, and overcoat are also included in the model. The material properties and thickness of each layer are shown in Table 1. These values, in particular the thermal conductivity, are process-and-thickness dependent, and we used the data published in various papers [10,14,21,22]. The thermal conductivities of thin layers are higher than the bulk values due to the heat carrier-boundary scattering and the altered microstructure of thin films [21]. A series of thermal-structural-coupled-field finite element analyses have been carried out using ANSYS, a commercial finite element package, to study the actuation performance of the thermal nanoactuator. The air-bearing modeling was done using the CML Air Bearing Simulator, which solved the generalized Reynolds equation to obtain the steady-state flying attitude.

At the head-disk interface heat is transferred from the slider to the disk through the air-bearing cooling effect. The cooling effect of the air bearing plays a key role in this 3D heat transfer problem. Chen et al. [23] found that the dominant factor of this effect was heat conduction. They applied the slip condition for the velocity and the jump condition for the temperature at the boundaries of the air bearing and obtained the heat transfer model as follows:

\[
q(x,y) = -k \frac{T_s(x,y) - T_d}{h(x,y) + 2\beta \lambda(x,y)} + f(x,y)
\]

where \(T_s\) and \(T_d\) are the temperatures of the slider and the disk, respectively, \(\beta = 2(1 - \sigma_T) / (\gamma + 1)\) \(Pr\), \(h\) is the FH (air-bearing thickness) of the slider, the mean free path of the air under pressure \(p\) is \(\lambda = \lambda_0 p / p\), and the thermal conductivity \(k\) is a very weak function of pressure.

Note that \(Ts\), \(h\), and \(\lambda\) are functions of coordinates \(x\) and \(y\). For air at \(T = 300\) K and atmospheric pressure, the thermal conductivity \(k = 0.0263\) \(W/mK\), the mean free path \(\lambda_0 = 65\) nm, and the Prandtl number \(Pr = 0.7\). The specific heat ratio \(\gamma\) is 1.4, and the thermal accommodation coefficient \(\sigma_T\) is 0.9. We assume that the disk surface is kept at the ambient temperature \(T_d = 25\) °C. Since the first term on the right-hand side of (1) is about 1–2 orders of magnitude larger than the other terms \(f(.)\), only this term is modeled in this study. Unlike the ABS, the dominant factor of heat

<table>
<thead>
<tr>
<th>Layer and material</th>
<th>Young’s modulus (GPa)</th>
<th>Thermal conductivity (W/m K)</th>
<th>Coefficient of thermal expansion ((×10^{-6}/°C))</th>
<th>Specific heat (J/kg K)</th>
<th>Poisson’s ratio</th>
</tr>
</thead>
<tbody>
<tr>
<td>Slider substrate</td>
<td>(\text{Al}_2\text{O}_3)-TiC</td>
<td>380</td>
<td>20</td>
<td>7.9</td>
<td>878</td>
</tr>
<tr>
<td>Undercoat (1.2 μm)</td>
<td>(\text{Al}_2\text{O}_3)</td>
<td>200</td>
<td>1.5</td>
<td>7.5</td>
<td>760</td>
</tr>
<tr>
<td>Shields (2.0 μm)</td>
<td>Ni-Fe</td>
<td>207</td>
<td>35</td>
<td>12.2</td>
<td>470</td>
</tr>
<tr>
<td>Bottom pole (1.0 μm)</td>
<td>Ni-Fe</td>
<td>207</td>
<td>35</td>
<td>12.2</td>
<td>470</td>
</tr>
<tr>
<td>Coil (2. μm)</td>
<td>Cu</td>
<td>120</td>
<td>395</td>
<td>16.5</td>
<td>390</td>
</tr>
<tr>
<td>Heater (250 nm)</td>
<td>Ni-Fe (thin layer)</td>
<td>207</td>
<td>30</td>
<td>12.2</td>
<td>470</td>
</tr>
<tr>
<td>Coil insulation (5 μm)</td>
<td>Photoresist</td>
<td>7</td>
<td>0.19</td>
<td>51.0</td>
<td>1460</td>
</tr>
<tr>
<td>Top pole (1.0 μm)</td>
<td>Ni-Fe</td>
<td>207</td>
<td>35</td>
<td>12.2</td>
<td>470</td>
</tr>
<tr>
<td>Overcoat (39.7 μm)</td>
<td>(\text{Al}_2\text{O}_3)</td>
<td>200</td>
<td>1.5</td>
<td>7.5</td>
<td>760</td>
</tr>
</tbody>
</table>

Table 1 Material properties used in the FEA [10,14,21,22]
transfer at the non-ABS surfaces of the slider is heat convection with a coefficient on the order of 100 W/m² K.

The numerical iteration approach developed in [20] is adopted in this paper. We first used the CML Air Bearing Simulator to obtain the nominal FH, pitch, roll and air pressure distribution of an ABS slider. Then we used Eq. (1) to specify the thermal boundary condition at the ABS and calculated the temperature distribution of the slider body. The obtained temperature distribution was then used as the body load to solve the slider deformation and actuated pole-tip protrusion due to the temperature gradient and the mismatch of coefficient of thermal expansions of various materials. Since the thermal protrusion causes deformation of the ABS and hence changes the flying attitudes and the thermal boundary conditions, several iterations are required to achieve an equilibrium solution.

3 Results and Discussions

3.1 Four Air-Bearing Designs. We study the actuation performances of four ABS designs that have distinct flying attitudes and pressure distributions, as shown in Fig. 2. Their flying atti-

<table>
<thead>
<tr>
<th>ABS Design</th>
<th>FH (nm)</th>
<th>Pitch (µrad)</th>
<th>Roll (µrad)</th>
<th>Peak pressure (atm)</th>
<th>Linear velocity (m/s)</th>
<th>Skew (deg)</th>
</tr>
</thead>
<tbody>
<tr>
<td>CML-5nm</td>
<td>5.3</td>
<td>220</td>
<td>0.9</td>
<td>18.9</td>
<td>17.3</td>
<td>9.10</td>
</tr>
<tr>
<td>Slider A</td>
<td>11.5</td>
<td>130</td>
<td>−1.2</td>
<td>23.4</td>
<td>37.5</td>
<td>−2.56</td>
</tr>
<tr>
<td>Slider B</td>
<td>12.0</td>
<td>115</td>
<td>2.4</td>
<td>21.0</td>
<td>18.0</td>
<td>−2.56</td>
</tr>
<tr>
<td>Scorpion III</td>
<td>10.5</td>
<td>124</td>
<td>−0.4</td>
<td>38.0</td>
<td>37.5</td>
<td>−2.56</td>
</tr>
</tbody>
</table>

Fig. 2 Four ABS designs used in this study (length: 1.25 mm; width: 1.00 mm). Different colors indicate different etching levels.

Fig. 3 Air pressure distributions of the ABS sliders. The scale displayed is normalized to ambient pressure: \((p−p_a)/p_p\).
Fig. 4 Prototype of a fabricated Al2O3-TiC slider having an isolated center trailing pad with all relevant read/write elements. The dimension of the pad is 100 μm by 30 μm, which can also be increased to accommodate larger transducer designs.

Fig. 5 (a) Simulation of gap FH and minimum FH profiles of Scorpion III at sea level, 0 m, and high altitude, 4500 m; and (b) simulation of pitch and roll profiles of Scorpion III at sea level, 0 m. The skew angle varies from −15.62 deg to 7.22 deg from the inner diameter to the outer diameter.

3.2 Steady-State Analysis. Figure 6 shows the steady-state heat transfer film coefficients on the air-bearing surfaces of the four designs obtained after several numerical iterations at a heating power of 20 mW. Only the part of the ABS that is close to the trailing edge is plotted. The distances of the write gap and giant magnetoresistive (GMR) sensor from the trailing edge are 33 μm and 36.5 μm, respectively. The disk linear speeds and skew angles are given in Table 2. It is seen that the interface thermal conductance is not constant and is, indeed, a strong function of both the FH and air pressure distributions. The peak values are about 1.6, 1.2, 1.4, and 0.2 MW/m²K for CML-5 nm, slider A, slider B, and Scorpion, respectively. The value of the Scorpion ABS is ~83–86% less than those of sliders A and B, even though the FHs of the three designs are similar, which clearly indicates that the effect of the pressure is more significant than the FH effect.

Figures 7 and 8 show the comparisons of the temperature rise and heat flux of the four designs at 20 mW, respectively. The maximum temperature rises are 2.78 K, 2.52 K, 3.21 K, and 7.8 K for CML-5 nm, slider A, slider B, and Scorpion, respectively. The Scorpion ABS has a higher temperature increase due to its relatively low heat flux, whereas the temperature distribution of slider B exhibits a different pattern from the others.

Using the temperature distributions as body loads in the finite element models, we carried out the static structural analysis to calculate the slider deformation. A comparison of the protrusion profiles is shown in Fig. 9. The maximum A-PTPs on the ABSs are found to be 3.86 nm, 3.56 nm, 4.54 nm, and 6.86 nm for CML-5 nm, slider A, slider B, and Scorpion, respectively. As expected, the Scorpion slider achieves 51% more protrusion than the second highest one at the same heating power of 20 mW. Figure 10 shows the protrusion profiles along the center line across the read/write elements, which indicates that the peaks of protrusions are located at the read/write elements as a result of the higher local temperature and higher coefficients of thermal expansion of the metal layers.

Figure 11 shows the FH reductions as a function of the dc heating power. The results indicate that the FH reduction is not proportional to the power. Instead, quadratic expressions have been found to best fit the curves for all the ABSs. This nonlinearity is related to the fact that the air bearing becomes stiffer when the FH is reduced, and the heat transfer across the air bearing also
Fig. 6  Distributions of the interface thermal conductance on the air-bearing surfaces at a heating power of 20 mW. Only part of the ABS that is close to the trailing edge is plotted. The distances of the write gap and the GMR sensor from the trailing edge are 33 μm and 36.5 μm, respectively.

Fig. 7  Distributions of the temperature rises on the ABSs at a heating power of 20 mW
Fig. 8 Distributions of the heat flux on the ABSs at a heating power of 20 mW

Fig. 9 Distributions of the A-PTP on the ABSs at a heating power of 20 mW
becomes more effective because of the increased pressure and reduced FH.

In order to evaluate the performance of thermal actuation, we defined five measures of merit as follows:

1. Actuation efficiency (in percent): The ratio of FH reduction to A-PTP
2. Power consumption (in megawatts): The power required for lowering one unit of the FH
3. Peak pressure increase (in atmospheres): The increase of pressure peak caused by the thermal protrusion
4. Protrusion rate (in nanometers per megawatt): The amount of protrusion per unit power
5. Temperature rise of the sensor (in Kelvin): The temperature rise of the GMR sensors as a function of FH

Figure 12 shows a comparison of the A-PTP as a function of the heating power. Scorpion exhibits an increase of 80%, 104%, and 63% in the protrusion rate over CML-5 nm, slider A and slider B, respectively. Figure 13 shows the actuation efficiency as a function of the FH for the four ABSs. It is noted that the efficiencies of CML-5 nm, sliders A and B monotonically decrease as the FHs are reduced by the thermal protrusions and the values range from 40% to 60%. However, Scorpion demonstrates virtually 100% efficiency, which does not depend on the FH. The heating powers required for lowering the FHs are shown in Fig. 14. The Scorpion ABS requires remarkably less power for reducing the FH from 10 nm to 3 nm compared to the other designs. Another important parameter that has to be considered is the pressure increase due to the deformed ABS. Figure 15 shows that the peak pressures of...
sliders A and B are 84 atm and 149 atm, respectively, when the FHs are reduced to 3 nm. Such high pressures may not be physical in reality and may cause adverse effects, such as slider deformation. Since the Scorpion slider is supported by the two side pressure peaks and the pressure underneath the center pad is relatively low, the thermal protrusion of the center pad does not affect the pressure distribution.

The temperature rise of the read/write elements is of great concern in the thermal nanoactuator. The read-back signal of GMR sensors can be significantly altered by thermal influences since their electrical resistance is temperature dependent. Figure 16 shows the temperature rises of the sensors as a function of the FH. It is observed that Scorpion has less temperature rise at FHs over 5 nm compared to sliders A and B. The temperatures of sliders A and B decrease when the FHs are reduced to <5 nm due to the highly concentrated pressures.

3.3 Transient Analysis. The bandwidth of thermal actuation is of great importance because it determines the response time of the thermal protrusion to the heating power. A transient thermal study was conducted to investigate the bandwidth of the thermal protrusion when the slider flies over a disk. The power required for the first 1 nm FH reduction for each of the ABSs was applied from 0 to 2.5 ms and was turned off at 2.5 ms. The temperature changes of the GMR sensors were monitored as shown in Fig. 17. It requires about 1–2 ms for the read/write transducers to reach their steady-state values, corresponding to a bandwidth of 0.5–1 kHz. It is also seen that Scorpion has the least temperature rise and takes longer to reach its steady-state value, which implies a trade-off between increasing the bandwidth and decreasing the temperature rise of the sensors.

4 Conclusions

The effects of ABS on thermal actuation have been studied by numerical simulation. A series of three-dimensional thermal-structural finite element analyses were conducted using velocity slip and temperature jump boundary conditions to formulate the heat transfer across the head-disk interface. An iteration procedure was used to obtain the equilibrium solutions. Four ABS designs with distinct features were simulated. In order to evaluate the performance of thermal actuation, we defined five measures of merit, including protrusion rate, actuation efficiency, power consumption, pressure peak, and temperature rise of the sensor. We found that the efficiencies of three conventional designs decrease as the FHs are continuously reduced. A new slider, Scorpion, which meets all design and fabrication requirements, has been presented and exhibits an overall enhancement, including virtually 100% efficiency with significantly less power consumption. Quadratic expressions have been found to best fit the curves of the FH reduction as a function of the heating power for all the designs.

Transient thermal analysis of the sliders in the flying conditions with a varying heating power showed that about 1–2 ms are required for the temperature to reach the steady-state values for the first nanometer actuation. It is found that Scorpion has the least temperature rise of the GMR sensor at the first 1 nm FH reduction, but the response time was longer than the other three designs. Therefore, there is a trade-off between increasing the actuation bandwidth and decreasing the power consumption.

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References


Fig. 17 Transient temperature changes of the flying sliders with a varying heating power. The power required for the first 1 nm FH reduction for each of the ABSs was applied from 0 to 2.5 ms and was turned off at 2.5 ms.