# Chapter 12: Boundary-Value Problems in Rectangular Coordinates

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December 18, 2013

In this lecture, we focus on solving some *classical* partial differential equations in boundary-value problems.

Instead of solving the general solutions, we are only interested in finding useful particular solutions.

We focus on linear second order PDE:  $(A, \dots, G)$ : functions of x, y

$$Au_{xx} + Bu_{xy} + Cu_{yy} + Du_x + Eu_y + Fu = G.$$

notation: for example, 
$$u_{xy} := \frac{\partial^2 u}{\partial x \partial y}$$
.

**Method**: Separation of variables – convert a PDE into two ODE's

#### Types of Equations:

- Heat Equation
- Wave Equation
- Laplace Equation

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### Classification of Linear Second Order PDE

$$Au_{xx} + Bu_{xy} + Cu_{yy} + Du_x + Eu_y + Fu = G.$$

notation: for example, 
$$u_{xy} := \frac{\partial^2 u}{\partial x \partial y}$$
.

1 Homogeneous vs. Nonhomogeneous

Homogeneous 
$$\iff G = 0$$
  
Nonhomogeneous  $\iff G \neq 0$ .

**2** Hyperbolic, Parabolic, and Elliptic:  $A, B, C, \dots, G$ : constants,

Hyperbolic	$\iff B^2 - 4AC > 0$
Parabolic	$\iff B^2 - 4AC = 0$
Elliptic	$\iff B^2 - 4AC < 0$

## Superposition Principle

#### **Theorem**

If  $u_1(x, y), u_2(x, y), \dots, u_k(x, y)$  are solutions of a homogeneous linear PDE, then a linear combination

$$u(x,y) := \sum_{n=1}^{k} c_n u_n(x,y)$$

is also a solution.

**Note**: We shall assume without rigorous argument that the linear combination can be an infinite series

$$u(x,y) := \sum_{n=1}^{\infty} c_n u_n(x,y)$$

1 Separation of Variables and Classical PDE's

2 Wave Equation

3 Laplace's Equation

# Separation of Variables

To find a particular solution of an PDE, one method is **separation of variables**, that is, assume that the solution u(x, y) takes the form of a product of a x-function and a y-function:

$$u(x, y) = X(x) Y(y).$$

Then, with the following, *sometimes* the PDE can be converted into an ODE of X and an ODE of Y:

$$\begin{split} u_x &= \frac{dX}{dx}Y = X'Y, & u_y &= X\frac{dY}{dy} = XY' \\ u_{xx} &= \frac{d^2X}{dx^2}Y = X''Y, & u_{yy} &= X\frac{d^2Y}{dy^2} = XY'', & u_{xy} &= X'Y' \end{split}$$

**Note**: Derivatives are with respect to different independent variables. For example,  $X':=\frac{dX}{dx}$ .

### Convert a PDE into Two ODE's

#### Example

Use separation of variables to convert the PDE below into two ODE's.

$$x^{2}u_{xx} + (x+1)u_{y} + (x+xy)u = 0$$

With u(x, y) = X(x) Y(y), the original PDE becomes

$$x^{2}X''Y + (x+1)XY' + (x+1)yXY = 0$$

$$\implies x^{2}X''Y = -(x+1)X(Y' + yY)$$

$$\implies \frac{x^{2}X''}{(x+1)X} = -\frac{Y'}{Y} - y = \lambda \quad \text{separation constant}$$

Left-hand side is a function of x, independent of y, Right-hand side is a function of y, independent of x. Hence, the above is equal to something independent of x and y

### Convert a PDE into Two ODE's

#### Example

Use separation of variables to convert the PDE below into two ODE's.

$$x^{2}u_{xx} + (x+1)u_{y} + (x+xy)u = 0$$

With u(x, y) = X(x) Y(y), the original PDE becomes

$$\begin{split} x^2X''Y + (x+1)XY' + (x+1)yXY &= 0 \\ \Longrightarrow x^2X''Y &= -(x+1)X(Y'+yY) \\ \Longrightarrow \frac{x^2X''}{(x+1)X} &= -\frac{Y'}{Y} - y = \lambda \quad \text{separation constant} \\ \Longrightarrow \begin{cases} x^2X''(x) - \lambda(x+1)X(x) &= 0 \\ Y'(y) + (y+\lambda)Y(y) &= 0 \end{cases} \end{split}$$

### Some Remarks

- 1 The method of separation of variables can only solve for *some* linear second order PDE's, not all of them.
- 2 For the PDE's considered in this lecture, the method works.
- $\blacksquare$  The method may work for both homogeneous (G=0) and nonhomogeneous ( $G \neq 0$ ) PDE's

$$Au_{xx} + Bu_{xy} + Cu_{yy} + Du_x + Eu_y + Fu = G.$$

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### Three Classical PDE's

In this lecture we focus on solving boundary-value problems of the following three classical PDE's that arise frequently in physics, mechanics, and engineering:

1 (One-Dimensional) Heat Equation/Diffusion Equation

$$ku_{xx} = u_t, \ k > 0$$

2 (One-Dimensional) Wave Equation/Telegraph Equation

$$a^2 u_{xx} = u_{tt}$$

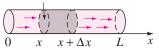
3 (Two-Dimensional) Laplace Equation

$$u_{xx} + u_{yy} = 0$$

## Heat Transfer within a Thin Rod: Heat Equation

### Assumptions:

Cross section of area A



- Heat only flows in x-direction.
- No heat escapes from the surface.
- No heat is generated in the rod.
- **Rod** is homogeneous with density  $\rho$ .

Let u(x, t) denote the temperature of the rod at location x at time t.

$$dQ = \gamma (\rho A dx) u \implies Q_x = \gamma \rho A u \implies Q_{xt} = \gamma \rho A u_t$$

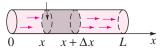
Heat transfer rate through the cross section  $= -KAu_x$ , and hence the net heat rate inside [x, x + dx] is  $dQ_t = -KAu_x(x, t) - (-KAu_x(x + dx, t))$ 

$$dQ_t = KAu(u_x(x+dx,t) - u_x(x,t)) = KAu_{xx}dx$$

$$\implies Q_{tx} = KAu_{xx}$$

## Heat Transfer within a Thin Rod: Heat Equation

#### Cross section of area A



#### Assumptions:

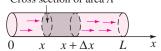
- Heat only flows in x-direction.
- No heat escapes from the surface.
- No heat is generated in the rod.
- **Rod** is homogeneous with density  $\rho$ .

Let u(x,t) denote the temperature of the rod at location x at time t. Hence,

$$\begin{cases} Q_{xt} = \gamma \rho A u_t \\ Q_{tx} = KA u_{xx} \end{cases} \implies \gamma \rho A u_t = KA u_{xx} \implies \left(\frac{K}{\gamma \rho}\right) u_{xx} = u_t$$
$$\implies \boxed{k u_{xx} = u_t, \ k > 0}$$

## Heat Equation: Initial and Boundary Conditions

#### Cross section of area A



#### Initial Condition:

Provides the spatial distribution of the temperature at time t=0.

$$u(x,0) = f(x), \ 0 < x < L$$

#### **Boundary Conditions:**

At the end points x = 0 and x = L, give the constraints on

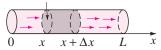
- ullet u: (Dirchlet condition), for example, ( $u_0$ : constant)  $u(L,t)=u_0 \quad \text{Temperature at the right end is held at constant.}$
- $u_x$ : (Neumann condition), for example,

$$u_x(L,t) = 0$$
 The right end is insulated.

 $u_x+hu$ : (Robin condition), for example,  $(h>0,u_m)$ : constants)  $u_x(L,t)=-h\left\{u(L,t)-u_m\right\} \quad \text{Heat is lost from the right end.}$ 

## Heat Equation: Boundary-Value Problems

#### Cross section of area A

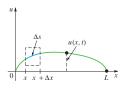


- A problem involving both initial and boundary conditions is called a boundary-value problem
- At the two boundaries x 0 and x = L, one can use different kinds of conditions.

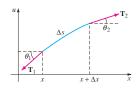
#### Examples:

$$\begin{array}{lll} ku_{xx} = u_t, & 0 < x < L, & t > 0 & & \text{Heat} \\ u(0,t) = u_0, & u_x(L,t) = -h\left\{u(L,t) - u_m\right\}, & t > 0 & \text{Boundary} \\ u(x,0) = f(x), & 0 < x < L & & \text{Initial} \end{array}$$

# Dynamics of a String Fixed at Two Ends: Wave Equation



(a) Segment of string



#### **Assumptions:**

- No external force.
- Tension force is large compared to gravity and is the same at all points.
- Slope of the curve is very small at all points.
- Vertical displacement ≪ string length.
- String has mass per unit length  $\rho$ .

Let u(x,t) denote the vertical position (displacement) of the string at location x at time t.

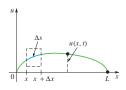
Consider the string in [x, x + dx]. Net vertical force is

$$T(\sin \theta_2 - \sin \theta_1) \approx T(\tan \theta_2 - \tan \theta_1)$$

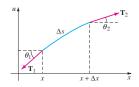
$$= T\{u_x(x + dx, t) - u_x(x, t)\}$$

$$= Tu_{xx}dx$$

# Dynamics of a String Fixed at Two Ends: Wave Equation



(a) Segment of string



#### Assumptions:

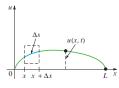
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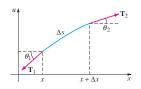
Since the slope is small, the mass  $\approx \rho dx$ . Hence

$$Tu_{xx} dx = (
ho dx) \ u_{tt} \implies rac{T}{
ho} u_{xx} = u_{tt}$$
  $\implies \boxed{a^2 u_{xx} = u_{tt}}$ 

# Wave Equation: Initial and Boundary Conditions



(a) Segment of string



#### Initial Conditions:

Provide the initial displacement u and velocity  $u_t$  at time t=0.

$$u(x,0) = f(x), u_t(x,0) = g(x), 0 < x < L$$

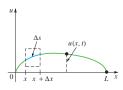
#### **Boundary Conditions:**

At the end points x=0 and x=L, give the constraints on u,  $u_x$ , or  $u_x+hu$ . Usually in the scenario of strings, the boundary conditions are

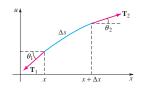
$$u(0,t)=0,\ u(0,L)=0,\ t>0$$
 Both ends are fixed.

$$u_x(0,t)=0,\ u_x(0,L)=0,\ t>0$$
 Free-ends condition

# Wave Equation: Boundary-Value Problems



(a) Segment of string



#### Examples:

#### Both ends are fixed:

$$a^{2}u_{xx} = u_{tt}, \quad 0 < x < L, \quad t > 0$$
  
 $u(0, t) = 0, \quad u(L, t) = 0, \quad t > 0$   
 $u(x, 0) = f(x), \quad u_{t}(x, 0) = g(x), \quad 0 < x < L$ 

equation Boundary

Wave

condition Initial condition

#### Free Ends:

$$a^{2} u_{xx} = u_{tt}, \quad 0 < x < L, \quad t > 0$$
  

$$u_{x}(0, t) = 0, \quad u_{x}(L, t) = 0, \quad t > 0$$
  

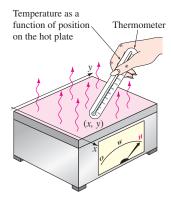
$$u(x, 0) = f(x), \quad u_{t}(x, 0) = g(x), \quad 0 < x < L$$

Wave equation Boundary condition

Initial condition

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## Laplace's Equation



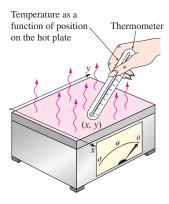
- Laplace's equation usually occurs in time-independent problems involving potentials.
- Its solution can also be interpreted as a steady-state temperature distribution.
- Two-dimensional Laplace Equation

$$\nabla^2 u := u_{xx} + u_{yy} = 0$$

■ Three-dimensional Laplace Equation

$$\nabla^2 u := u_{xx} + u_{yy} + u_{zz} = 0$$

# Laplace's Equation: Boundary Conditions



#### **Boundary Conditions:**

In the *x*-direction, at the end points x = 0 and x = a, give the constraints on u,  $u_x$ , or  $u_x + hu$ .

In the *y*-direction, at the end points y=0 and y=b, give the constraints on u,  $u_y$ , or  $u_y+hu$ .

#### Examples:

Both ends in x are insulated

$$u_x(0, y) = 0, \quad u_x(a, y) = 0$$

■ Temperatures of two ends in y are held at different distributions

$$u(x, 0) = f(x), \quad u(x, b) = g(x)$$

# Laplace's Equation: Boundary-Value Problems

#### Example:

$$\begin{array}{lll} u_{xx} + u_{yy} = 0, & 0 < x < a, & 0 < y < b & & \text{Laplace's equation} \\ u_x(0,y) = 0, & u_x(a,y) = 0, & 0 < y < b & & \text{Boundary condition} \\ u(x,0) = f(x), & u(x,b) = g(x), & 0 < x < a & & \text{Boundary condition} \end{array}$$

## Modifications of Heat and Wave Equations

In the derivation of the heat equation and the wave equation, we assume that there is no internal or external influences. For example, *no heat escapes from the surface, no heat is generated in the rod, no external force act on the string, etc.* 

Taking external and internal influences into account, more general forms of the heat equation and the wave equation are the following:

$$ku_{xx}+G(x,t,u,u_x)=u_t$$
 Heat Equation  $a^2u_{xx}+F(x,t,u,u_t)=u_{tt}$  Wave Equation

#### Example:

$$ku_{xx}-h(u-u_m)=u_t$$
 heat transfers from the surface to an environment with constant temperature  $u_m$  External force  $f$  acts on the string

# Homogeneous vs. Nonhomogeneous Boundary Conditions

#### **Homogeneous Boundary Condition:**

$$u_x(0, y) = 0$$
,  $u_x(a, y) = 0$ ,  $u(x, 0) = 0$ ,  $u(0, L) = 0$ 

#### Nonhomogeneous Boundary Condition:

$$u_x(0, y) = f(y), \quad u_x(a, y) = g(y), \quad u(x, L) = u_m$$

Typically, when using separation of variables, start with the independent variable associated with homogeneous boundary conditions.

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1 Separation of Variables and Classical PDE's

2 Wave Equation

3 Laplace's Equation

### Wave Equation: a Boundary-Value Problem

$$\begin{split} \text{Solve } u(x,t): \quad & au_{xx} = u_{tt}, \quad 0 < x < L, \quad t > 0 \\ \text{subject to:} \quad & u(0,t) = 0, \quad u(L,t) = 0, \quad t > 0 \\ & u(x,0) = f(x), \quad u_t(x,0) = g(x), \quad 0 < x < L \end{split}$$

We focus on solving the above BVP (both ends are fixed).

#### Step 1: Separation of variables:

Assume that the solution u(x,t) = X(x) T(t),  $X, T \neq 0$ . Then,

$$a^{2}u_{xx} = u_{tt} \implies a^{2}X''T = XT'' \implies \frac{X''}{X} = \frac{T''}{a^{2}T} = -\lambda$$

$$\implies \begin{cases} X'' + \lambda X &= 0\\ T'' + a^{2}\lambda T &= 0 \end{cases}$$

The 2 homogeneous boundary conditions become X(0) = X(L) = 0.

### Solve in the x-Dimension and Find $\lambda$

Solve : 
$$X'' + \lambda X = 0, \quad T'' + a^2 \lambda T = 0$$
 subject to : 
$$X(0) = 0, \quad X(L) = 0$$
 
$$X(x)T(0) = f(x), \quad X(x)T'(0) = g(x), \quad 0 < x < L$$

**Step 2**:  $\lambda$  remains to be determined. What values should  $\lambda$  take?

1 
$$\lambda = 0$$
:  $X(x) = c_1 + c_2 x$ .  $X(0) = X(L) = 0 \implies c_1 = c_2 = 0$ .

2 
$$\lambda = -\alpha^2 < 0$$
:  $X(x) = c_1 e^{-\alpha x} + c_2 e^{\alpha x}$ .  
Plug in  $X(0) = X(L) = 0$ , we get  $c_1 = c_2 = 0$ .

3  $\lambda=\alpha^2>0$ :  $X(x)=c_1\cos(\alpha x)+c_2\sin(\alpha x)$ . Plug in X(0)=X(L)=0, we get  $c_1=0$ , and  $c_2\sin(\alpha L)=0$ . Hence,  $c_2\neq 0$  only if  $\alpha L=n\pi$ .

Since 
$$X \neq 0$$
, pick  $\lambda = \frac{n^2\pi^2}{L^2}$ ,  $n = 1, 2, \dots \Longrightarrow X(x) = c_2 \sin \frac{n\pi}{L} x$ .

### Solve in *t*-Dimension and Superposition

$$\begin{split} \text{Solve}: \quad X'' + \lambda X &= 0, \quad T'' + a^2 \lambda \, T = 0 \\ \text{subject to}: \quad X(0) &= 0, \quad X(L) = 0 \\ \quad X(x) \, T(0) &= \mathit{f}(x), \quad X(x) \, T'(0) = \mathit{g}(x), \quad 0 < x < L \end{split}$$

**Step 3**: Once we fix  $\lambda = \frac{n^2 \pi^2}{L^2}$ , n = 1, 2, ..., we obtain

$$X(x) = c_2 \sin\left(\frac{n\pi}{L}x\right), \quad T(t) = c_3 \cos\left(\frac{n\pi a}{L}t\right) + c_4 \sin\left(\frac{n\pi a}{L}t\right)$$

$$\implies u_n(x,t) = \left\{A_n \cos\left(\frac{n\pi a}{L}t\right) + B_n \sin\left(\frac{n\pi a}{L}t\right)\right\} \sin\left(\frac{n\pi}{L}x\right),$$

$$(A_n := c_2 c_3, \ B_n := c_2 c_4)$$

 $\implies u(x,t) := \sum_{n=1}^{\infty} u_n(x,t)$  is a solution, by the superposition principle.

## Plug in Initial Condition, Revoke Fourier Series, and Done

Solve 
$$u(x,t)$$
:  $au_{xx} = u_{tt}$ ,  $0 < x < L$ ,  $t > 0$  subject to:  $u(0,t) = 0$ ,  $u(L,t) = 0$ ,  $t > 0$  
$$u(x,0) = f(x)$$
,  $u_t(x,0) = g(x)$ ,  $0 < x < L$ 

**Step 4**: Plug in the initial conditions and find  $\{A_n, B_n \mid n = 1, 2, ...\}$ .

$$u(x,0) = f(x), \quad u(x,t) = \sum_{n=1}^{\infty} \left\{ A_n \cos\left(\frac{n\pi a}{L}t\right) + B_n \sin\left(\frac{n\pi a}{L}t\right) \right\} \sin\left(\frac{n\pi}{L}x\right)$$

$$\implies f(x) = \sum_{n=1}^{\infty} A_n \sin\left(\frac{n\pi}{L}x\right), \quad 0 < x < L$$

From the Fourier sine series expansion on (0, L), we get

$$A_n = \frac{2}{L} \int_0^L f(x) \sin\left(\frac{n\pi}{L}x\right) dx.$$

## Plug in Initial Condition, Revoke Fourier Series, and Done

Solve 
$$u(x,t): au_{xx} = u_{tt}, \quad 0 < x < L, \quad t > 0$$
 subject to :  $u(0,t) = 0, \quad u(L,t) = 0, \quad t > 0$  
$$u(x,0) = f(x), \quad u_t(x,0) = g(x), \quad 0 < x < L$$

**Step 4**: Plug in the initial conditions and find  $\{A_n, B_n \mid n = 1, 2, ...\}$ .

$$u_t(x,0) = g(x), \quad u(x,t) = \sum_{n=1}^{\infty} \left\{ A_n \cos\left(\frac{n\pi a}{L}t\right) + B_n \sin\left(\frac{n\pi a}{L}t\right) \right\} \sin\left(\frac{n\pi}{L}x\right)$$

$$\implies g(x) = \sum_{n=1}^{\infty} B_n \frac{n\pi a}{L} \sin\left(\frac{n\pi}{L}x\right), \quad 0 < x < L$$

From the Fourier sine series expansion on (0, L), we get

$$B_n \frac{n\pi a}{L} = \frac{2}{L} \int_0^L g(x) \sin\left(\frac{n\pi}{L}x\right) dx. \implies B_n = \frac{2}{n\pi a} \int_0^L g(x) \sin\left(\frac{n\pi}{L}x\right) dx.$$

### Final Solution

$$\begin{split} \text{Solve } u(x,t): \quad & au_{xx} = u_{tt}, \quad 0 < x < L, \quad t > 0 \\ \text{subject to}: \quad & u(0,t) = 0, \quad u(L,t) = 0, \quad t > 0 \\ & u(x,0) = f(x), \quad u_t(x,0) = g(x), \quad 0 < x < L \end{split}$$

#### Step 5: The final solution is

$$u(x,t) = \sum_{n=1}^{\infty} \left\{ A_n \cos\left(\frac{n\pi a}{L}t\right) + B_n \sin\left(\frac{n\pi a}{L}t\right) \right\} \sin\left(\frac{n\pi}{L}x\right)$$

$$= \sum_{n=1}^{\infty} C_n \sin\left(\frac{n\pi a}{L}t + \phi_n\right) \sin\left(\frac{n\pi}{L}x\right)$$

$$A_n = \frac{2}{L} \int_0^L f(x) \sin\left(\frac{n\pi}{L}x\right) dx, \quad B_n = \frac{2}{n\pi a} \int_0^L g(x) \sin\left(\frac{n\pi}{L}x\right) dx$$

$$C_n = \sqrt{A_n^2 + B_n^2}, \quad \sin\phi_n = \frac{A_n}{C_n}, \cos\phi_n = \frac{B_n}{C_n}$$

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### Standing Waves

The final solution

$$u(x,t) = \sum_{n=1}^{\infty} C_n \sin\left(\frac{n\pi a}{L}t + \phi_n\right) \sin\left(\frac{n\pi}{L}x\right)$$

is a linear combination of standing waves or normal modes

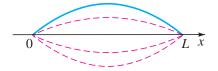
$$u_n(x,t) = C_n \sin\left(\frac{n\pi a}{L}t + \phi_n\right) \sin\left(\frac{n\pi}{L}x\right), \ n = 1, 2, \dots$$

For a normal mode n, at a fixed location x, the string moves with

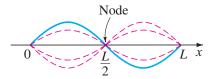
- time-varying amplitude  $C_n \sin\left(\frac{n\pi}{T}x\right)$
- frequency  $f_n := \frac{n\pi a/L}{2\pi} = \frac{na}{2L}$

Fundamental Frequency:  $f_1 := \frac{\pi a/L}{2\pi} = \frac{a}{2L}$ 

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(a) First standing wave



**(b)** Second standing wave

1 Separation of Variables and Classical PDE's

2 Wave Equation

3 Laplace's Equation

## Laplace's Equation: a Boundary-Value Problem

$$\begin{aligned} \text{Solve } u(x,y): \quad u_{xx} + u_{yy} &= 0, \quad 0 < x < a, \quad 0 < y < b \\ \text{subject to}: \quad u_x(0,y) &= 0, \quad u_x(a,y) &= 0, \quad 0 < y < b \\ u(x,0) &= 0, \quad u(x,b) &= f(x), \quad 0 < x < a \end{aligned}$$

We focus on solving the above BVP (both ends x = 0 and x = a are insulated).

#### **Step 1**: Separation of variables:

Assume that the solution u(x,y) = X(x) Y(y),  $X, Y \neq 0$ . Then,

$$u_{xx} + u_{yy} = 0 \implies X''Y + XY'' = 0 \implies \frac{X''}{X} = -\frac{Y''}{Y} = -\lambda$$

$$\implies \begin{cases} X'' + \lambda X &= 0 \\ Y'' - \lambda Y &= 0 \end{cases}$$

The 3 homogeneous boundary conditions become X'(0) = X'(a) = Y(0) = 0.

### Solve in the x-Dimension and Find $\lambda$

Solve : 
$$X'' + \lambda X = 0, \quad Y'' - \lambda \, Y = 0$$
 subject to : 
$$X'(0) = 0, \quad X'(a) = 0$$
 
$$Y(0) = 0, \quad X(x) \, Y(b) = \mathit{f}(x), \quad 0 < x < a$$

**Step 2**:  $\lambda$  remains to be determined. What values should  $\lambda$  take?

1 
$$\lambda = 0$$
:  $X(x) = c_1 + c_2 x$ .  $X'(0) = X'(a) = 0 \implies c_2 = 0$ .

2 
$$\lambda = -\alpha^2 < 0$$
:  $X(x) = c_1 e^{-\alpha x} + c_2 e^{\alpha x}$ .  
Plug in  $X'(0) = X'(a) = 0$ , we get  $c_1 = c_2 = 0$ .

3  $\lambda=\alpha^2>0$ :  $X(x)=c_1\cos(\alpha x)+c_2\sin(\alpha x)$ . Plug in X'(0)=X'(a)=0, we get  $c_2=0$ , and  $c_1\alpha\sin(\alpha a)=0$ . Hence,  $c_1\neq 0$  only if  $\alpha a=n\pi$ .

Since 
$$X \neq 0$$
, pick  $\lambda = \frac{n^2 \pi^2}{a^2}$ ,  $n = 0, 1, 2, \dots \Longrightarrow X(x) = c_1 \cos\left(\frac{n\pi}{a}x\right)$ .

## Solve in y-Dimension and Superposition

Solve : 
$$X'' + \lambda X = 0, \quad Y'' - \lambda \, Y = 0$$
 subject to : 
$$X'(0) = 0, \quad X'(a) = 0$$
 
$$Y(0) = 0, \quad X(x) \, Y(b) = \mathit{f}(x), \quad 0 < x < a$$

**Step 3**: Once we fix  $\lambda = \frac{n^2\pi^2}{\sigma^2}$ , n = 0, 1, 2, ..., we obtain  $X(x) = c_1 \cos\left(\frac{n\pi}{\sigma}x\right)$ 

$$Y(y) = \begin{cases} \mathscr{A} + c_4 y, & n = 0 \\ c_3 \cosh\left(\frac{n\pi}{a}y\right) + c_4 \sinh\left(\frac{n\pi}{a}y\right), & n \ge 1 \end{cases} \quad (Y(0) = 0 \implies c_3 = 0)$$

$$\implies u_n(x, y) = \begin{cases} A_0 y, & n = 0 \\ A_n \cos\left(\frac{n\pi}{a}x\right) \sinh\left(\frac{n\pi}{a}y\right), & n \ge 1 \end{cases}, \quad (A_n := c_1 c_4)$$

$$\implies u(x, y) := \sum_{n=0}^{\infty} u_n(x, y) \text{ is a solution, by the superposition principle.}$$

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# Plug in Initial Condition, Revoke Fourier Series, and Done

Solve 
$$u(x,y)$$
:  $u_{xx} + u_{yy} = 0$ ,  $0 < x < a$ ,  $0 < y < b$  subject to:  $u_x(0,y) = 0$ ,  $u_x(a,y) = 0$ ,  $0 < y < b$   $u(x,0) = 0$ ,  $u(x,b) = f(x)$ ,  $0 < x < a$ 

**Step 4**: Plug in the initial conditions and find  $\{A_n \mid n=1,2,\ldots\}$ .

$$u(x, b) = f(x), \quad u(x, y) = A_0 y + \sum_{n=1}^{\infty} A_n \cos\left(\frac{n\pi}{a}x\right) \sinh\left(\frac{n\pi}{a}y\right)$$

$$\implies f(x) = A_0 b + \sum_{n=1}^{\infty} A_n \cos\left(\frac{n\pi}{a}x\right) \sinh\left(\frac{n\pi}{a}b\right), \quad 0 < x < a$$

From the Fourier cosine series expansion on (0, a), we get

$$2A_0b = \frac{2}{a} \int_0^a f(x) dx, \quad A_n \sinh\left(\frac{n\pi}{a}b\right) = \frac{2}{a} \int_0^a f(x) \cos\left(\frac{n\pi}{a}x\right) dx$$

### Final Solution

#### Step 5: The final solution is

$$u(x,y) = A_0 y + \sum_{n=1}^{\infty} A_n \cos\left(\frac{n\pi}{a}x\right) \sinh\left(\frac{n\pi}{a}y\right)$$
$$A_0 = \frac{1}{ab} \int_0^a f(x) dx$$
$$A_n = \frac{2}{a \sinh\left(\frac{n\pi}{a}b\right)} \int_0^a f(x) \cos\left(\frac{n\pi}{a}x\right) dx, \quad n \ge 1$$

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