

INTRODUCTION TO SYMPLECTIC GEOMETRY

FOR NOVEMBER 4

1. ON THE COADJOINT ORBIT

Let G be a Lie group, \mathfrak{g} be its Lie algebra, and \mathfrak{g}^* be the dual vector space of \mathfrak{g} . The adjoint representation of G on \mathfrak{g} is defined by

$$\mathrm{Ad}_g(Y) = \left. \frac{d}{dt} \right|_{t=0} (g e^{tY} g^{-1}) . \quad (1.1)$$

The coadjoint action of G on \mathfrak{g}^* is characterized by

$$\langle \mathrm{Ad}_g^*(\xi), Y \rangle = \langle \xi, \mathrm{Ad}_{g^{-1}}(Y) \rangle \quad (1.2)$$

for any $Y \in \mathfrak{g}$. Here, $\langle \cdot, \cdot \rangle : \mathfrak{g}^* \times \mathfrak{g} \rightarrow \mathbb{R}$ is the dual pairing. The purpose of this note is to explain that there is a *canonical symplectic form* on the *orbits* of the coadjoint action. Part of the material here is taken from [CdS1, Homework 17].

1.1. The vector fields associated to the adjoint and coadjoint action. For any $X \in \mathfrak{g}$, it induces a vector field on \mathfrak{g} by the adjoint action:

$$\begin{aligned} \text{The vector field at } Y &= \left. \frac{d}{dt} \right|_{t=0} \mathrm{Ad}_{e^{tX}}(Y) \\ &= \left. \frac{d}{dt} \right|_{t=0} \left. \frac{d}{ds} \right|_{s=0} e^{tX} e^{sY} e^{-tX} \\ &= \left. \frac{d}{dt} \right|_{t=0} (e^{-tX})_*(Y|_{e^{tX}}) = [X, Y] . \end{aligned} \quad (1.3)$$

The computation relies on the following facts:

- If we think Y as a *left* invariant vector field, $Y|_g = \left. \frac{d}{ds} \right|_{s=0} g e^{sY}$.
- On G , the flow generated by X for time t is tantamount to the *right* multiplication by e^{tX} . This is basically the same as the previous fact.
- The equality (1.3) is an equivalent definition for the Lie derivative.

For any $X \in \mathfrak{g}^*$, denote by X^\sharp the vector field on \mathfrak{g}^* induced by the coadjoint action. Its value at ξ is characterized by

$$\begin{aligned} \langle X^\sharp|_\xi, Y \rangle &= \left. \frac{d}{dt} \right|_{t=0} \langle \mathrm{Ad}_{e^{tX}}^* \xi, Y \rangle \\ &= \left. \frac{d}{dt} \right|_{t=0} \langle \xi, \mathrm{Ad}_{e^{-tX}} Y \rangle = \langle \xi, [Y, X] \rangle . \end{aligned} \quad (1.4)$$

1.2. The skew-symmetric bilinear form. For any $\xi \in \mathfrak{g}^*$, define a skew-symmetric bilinear form on \mathfrak{g} by

$$\omega_\xi(X, Y) = \langle \xi, [X, Y] \rangle . \quad (1.5)$$

We now examine the kernel of this skew-symmetric bilinear form. If $X \in \ker(\omega_\xi)$, we apply (1.4) to find that

$$0 = \omega_\xi(X, Y) = -\langle X^\#|_\xi, Y \rangle$$

for any $Y \in \mathfrak{g}$. Therefore, $X^\#$ vanishes at ξ . It follows that $\ker(\omega_\xi)$ is the Lie algebra of the stabilizer of the action (1.2) at ξ .

We denote the stabilizer by G_ξ , and its Lie algebra by \mathfrak{g}_ξ . Since $\ker(\omega_\xi) = \mathfrak{g}_\xi$, ω_ξ induces a non-degenerate skew-symmetric bilinear form on $\mathfrak{g}/\mathfrak{g}_\xi$. This quotient space is identified with the *tangent space* of the *orbit* of the coadjoint action at ξ .

Hence, ω induces a non-degenerate 2-form on the orbit of the coadjoint action. It is defined by

$$\omega_\xi(X^\#|_\xi, Y^\#|_\xi) = \langle \xi, [X, Y] \rangle . \quad (1.6)$$

Notice that (1.5) is *not* a 2-form on \mathfrak{g}^* . The inputs are elements of \mathfrak{g} , but not \mathfrak{g}^* .

1.3. Exterior derivative of ω . We calculate the exterior derivative of (1.6). Due to the above comment, it only makes sense to compute the exterior derivative on the orbit.

$$\begin{aligned} (d\omega)(X^\#, Y^\#, Z^\#) &= X^\#(\omega(Y^\#, Z^\#)) + Z^\#(\omega(X^\#, Y^\#)) + Y^\#(\omega(Z^\#, X^\#)) \\ &\quad - \omega([Y^\#, Z^\#], X^\#) - \omega([X^\#, Y^\#], Z^\#) - \omega([Z^\#, X^\#], Y^\#) . \end{aligned}$$

For the terms of the first line,

$$\begin{aligned} X^\#(\omega(Y^\#, Z^\#)) &= X^\#(\langle \xi, [Y, Z] \rangle) \\ &= \left. \frac{d}{dt} \right|_{t=0} \langle \text{Ad}_{e^{tX}}^* \xi, [Y, Z] \rangle \\ &= \left. \frac{d}{dt} \right|_{t=0} \langle \xi, \text{Ad}_{e^{-tX}} [Y, Z] \rangle \\ &= -\langle \xi, [X, [Y, Z]] \rangle . \end{aligned}$$

For the terms of the second line,

$$\begin{aligned} \omega([Y^\#, Z^\#], X^\#) &= \omega([Y, Z]^\#, X^\#) \\ &= \langle \xi, [[Y, Z], X] \rangle . \end{aligned}$$

By the Jacobi identity, $d\omega = 0$. Thus, ω defines a symplectic form on the orbit of the coadjoint action. It is also known as the *Kostant–Kirillov symplectic structure*.

1.4. **The moment map.** Surely the Lie group G acts on the coadjoint orbit. It turns out that this group action is Hamiltonian, and the moment map is very simple.

For any $X \in \mathfrak{g}$, we claim that $\iota_{X^\#}\omega$ is d-exact. By (1.6) and (1.4),

$$\begin{aligned}\iota_{Y^\#}\iota_{X^\#}\omega &= \langle \xi, [X, Y] \rangle \\ &= \langle Y^\#|_\xi, X \rangle = Y^\#(\langle \xi, X \rangle)\end{aligned}$$

for any $Y \in \mathfrak{g}$. It follows that $\iota_{X^\#}\omega = d(\langle \xi, X \rangle)$.

Hence, we can just take the moment map to be ξ , and it is automatically equivariant with respect to the coadjoint action. In other words, the moment map of the coadjoint orbit is the *inclusion map* to \mathfrak{g}^* .

1.5. **Remarks.**

- In fact, we can apply this framework for $G = \text{SO}(3)$. Then we will get the answer for (6) of the Midterm.
- For Item (1.e) of Homework 8, the symplectic form should be

$$\omega_\xi(X^\#, Y^\#) = i \text{trace}([X, Y]\xi) .$$

The correct statement of the first fact should be

$$\begin{aligned}(d\omega)(X^\#, Y^\#, Z^\#) &= X^\#(\omega(Y^\#, Z^\#)) + Z^\#(\omega(X^\#, Y^\#)) + Y^\#(\omega(Z^\#, X^\#)) \\ &\quad - \omega([Y^\#, Z^\#], X^\#) - \omega([X^\#, Y^\#], Z^\#) - \omega([Z^\#, X^\#], Y^\#) .\end{aligned}$$

For the second fact, we consider the derivative of

$$\xi \mapsto \omega_\xi(X^\#, Y^\#) = i \text{trace}([X, Y]\xi)$$

along $Z^\#$, which is equal to

$$\frac{d}{dt}\Big|_{t=0} i \text{trace}([X, Y](e^{tZ}\xi e^{-tZ})) .$$

As mentioned above, the whole business of the symplectic form computation is done on the coadjoint orbit.