

**NONCOMMUTATIVITY  
NONASSOCIATIVITY  
AND NONLOCALITY  
IN STRING THEORY**

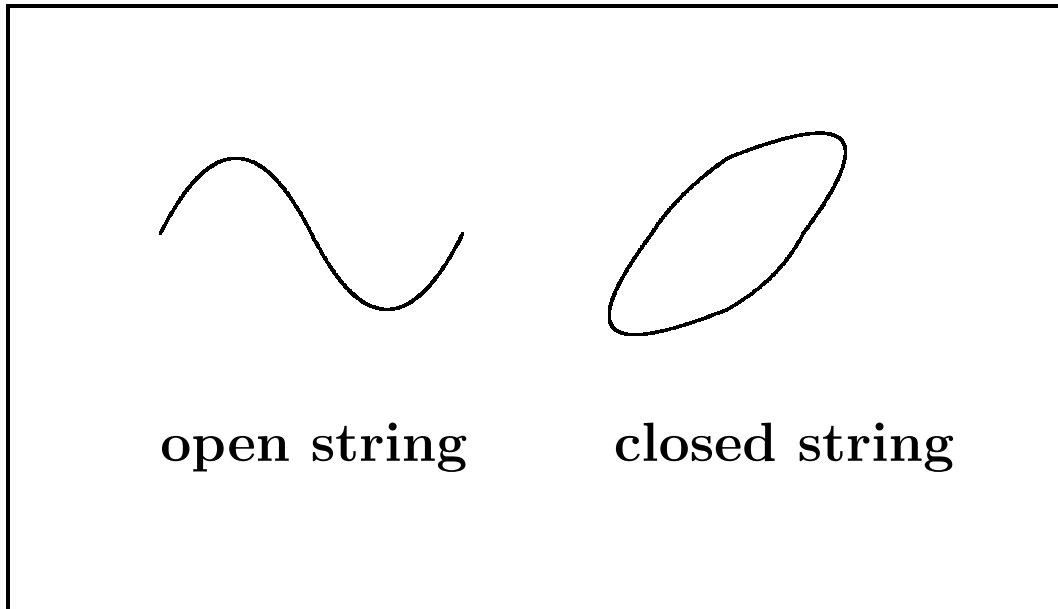
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Talk given at the symposium in celebration  
of Prof. C. N. Yang's 80th birthday.

String theory is a promising candidate for quantum gravity (and everything else).



String  $\simeq \infty$  particles

Nature of Space-Time

String Theory  $\rightarrow$  9+1 dimensions.

Geometry  $\leftrightarrow$  Physics

Euclidean Geometry

*(Newtonian Physics)*



Riemannian Geometry

*(General Relativity – Gravity)*



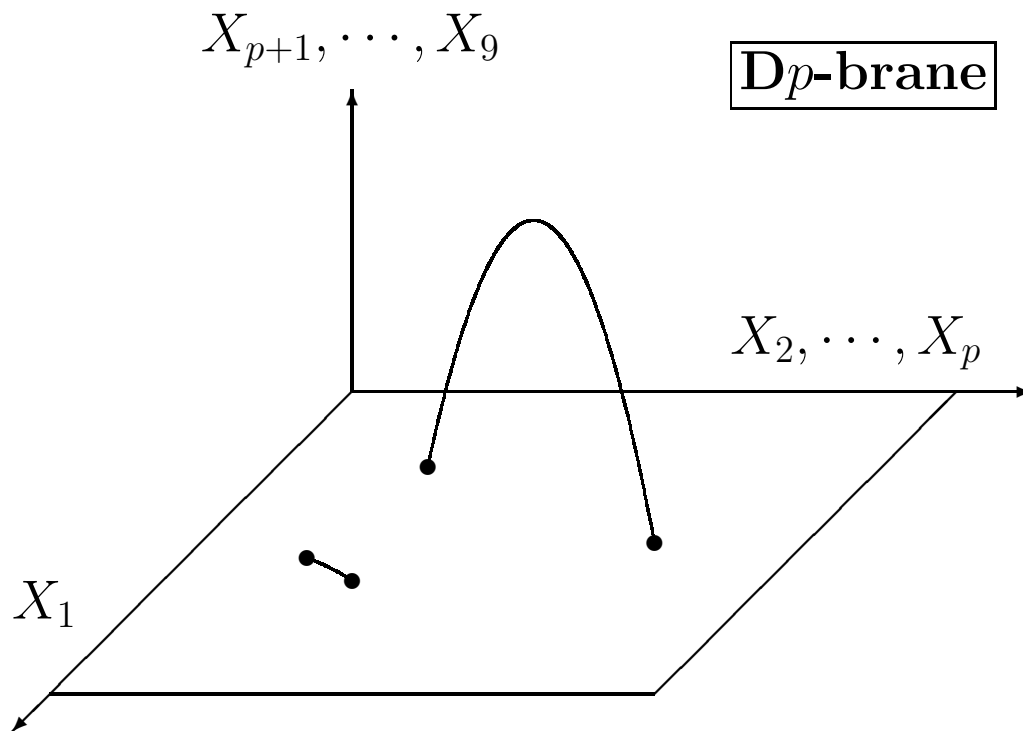
Non-Commutative Geometry

*(Quantum Gravity, String Theory)*

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Quantum Spacetime:

[Yang: On Quantized Spacetime, Phys. Rev. 72 (47) 874 ]



$Dp$ -brane  $\simeq$

$p$ -dim. subspace on which open strings end

Physics on  $N$  coincident D-branes  $\simeq$

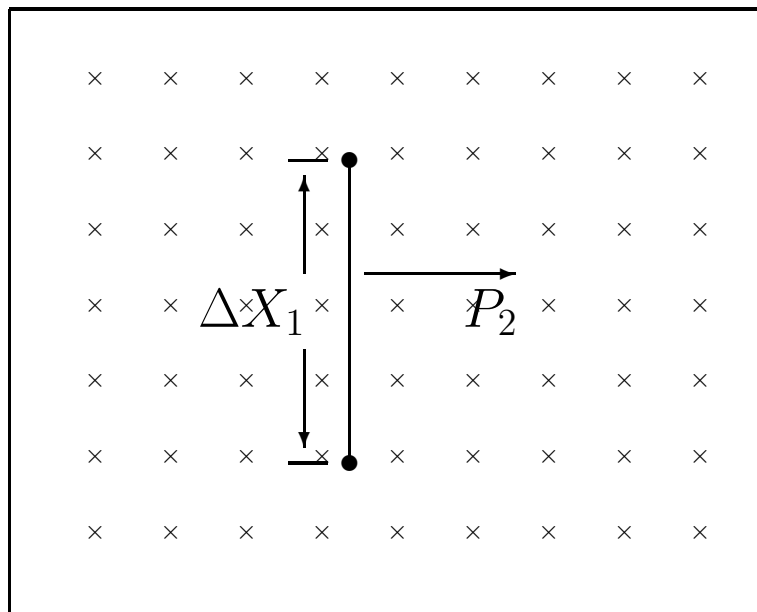
$SU(N)$  Yang-Mills theory

# D-brane with Background Field

**B field background:** *[Chu+Ho(98)]*

$$S = S_0 + \int B$$

**$B = 2$ -form gauge field**



**Nonlocality:**

$$\Delta X_1 = B_{12} P_2$$

$$\Rightarrow [X_1, X_2] \sim [\Delta X_1, X_2] = iB_{12}$$

## Dipole $\rightarrow$ Nonlocal $\rightarrow$ NC

Quantization of open string [*Chu+Ho(98,99)*]

$$\Rightarrow [X_i, X_j] = \left( \frac{B}{1 + B^2} \right)_{ij} \simeq iB_{ij}^{-1}$$

Noncommutativity effectively accounts for the interactions with the background  $B$  field.

Correlation functions for NC theory  $\simeq$   
Corr. fx's for  $B$  field background

[*Schomerus (99)*] [*Seiberg+Witten(99)*]

$H = dB = 0 \rightarrow$  Associativity

$$(fg)h = f(gh)$$

For  $H = dB \neq 0$ :

ASSOCIATIVE or NON-ASSOCIATIVE?

Correlation functions [*Cornalba+Shiappa(01)*]  $\rightarrow$

(indices omitted)

$$(f \bullet g) \bullet h - f \bullet (g \bullet h) = \frac{1}{6} B^{-1} B^{-1} B^{-1} H (\partial f) (\partial g) (\partial h) + \dots$$

open string quantization: [*Ho+Yeh(00)*]

$$[X, X] = iB^{-1} + \frac{1}{3} B^{-1} B^{-1} B^{-1} H P + \dots$$

$\Rightarrow$  associative algebra [*Ho(01)*]

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The algebra of  $X$  and  $P$  is mixed,  
like the quantized spacetime of Yang (47).

How to define gauge transformations?

[*Ho(01)*]

## Non-Commutative Spacetime

In the neighborhood of  $N$  D-branes,  
quantize a probing D-brane  $\Rightarrow$

$[H_0(00)]$

$$[X_i, X_j] = iL_{ij}$$

same as Yang (47) for 3+1 dim.

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For  $AdS_2 \times S^2$ ,  $[H_0 + L_i(00)]$

Noncommutativity  $\rightarrow$

uncertainty relation  $\rightarrow$

Bekenstein-Hawking entropy for charged black hole.



# NC → UNCERTAINTY → NONLOCALITY

## Fluctuations of background field

*[Chu+Ho+Kao(99)]*

→ Uncertainty principle of string theory

*[Yoneya(97), Li+Yoneya(98)]*

$$\Delta x \Delta t \geq 1/T_s$$

→ Nonlocality

a salient feature of string theory

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Perturbation theory of higher derivative  
or nonlocal theories

[Yang + Feldman:

The S-Matrix in the Heisenberg Representation,

Phys. Rev. 79, 972 (50)]

applied to string field theory *[Cheng+Ho+Yeh(01,02)]*

Nonlocality  $\sim$  limit of higher derivatives

Higher derivatives  $\rightarrow$   
more degrees of freedom

Nonlocality  $\rightarrow \infty$  degrees of freedom

*String*  $\simeq \infty$  *particles*

.....

String  $\simeq$  Nonlocal particle ?

Reparametrization symmetry of  
(*nonlocal*) particle worldline  $\simeq$   
Conformal symmetry of string worldsheet  
*[Ho(02)]*

## particle worldline Lagrangian

$$L = L(x, \dot{x}, \ddot{x}, \dots, x^{(n)})$$

$n = 1$ :

Phase space:  $x$  and its conjugate variable

Reparametrization:

$$\delta x = \epsilon \dot{x}$$

is generated by the Hamiltonian

$$[H, x] = i\dot{x}$$

$n = 2$ :

Phase space:  $x, \dot{x}$  and their conj. var.s

Reparametrization: (in addition to the one above)

$$\delta \dot{x} = (\epsilon \dot{x})' = \epsilon \ddot{x} + \dot{\epsilon} \dot{x}$$

$H$  only generates the term for  $\epsilon$ .

New generator needed for  $\dot{\epsilon}$ .

Theories with higher order time derivatives:

Phase space is larger.

Repara. has more independent parameters:

$$\epsilon, \dot{\epsilon}, \ddot{\epsilon}, \dots$$

→ More symmetry generators.

$$\begin{aligned}
L = & \\
& e \left[ A^{(0)}(x) + A_{\mu}^{(1)}(x) D x^{\mu} \right. \\
& + A_{\mu}^{(01)}(x) D^2 x^{\mu} + A_{\mu\nu}^{(20)}(x) D x^{\mu} D x^{\nu} \\
& + A_{\mu}^{(001)} D^3 x^{\mu} + A_{\mu\nu}^{(110)} D^2 x^{\mu} D x^{\nu} + A_{\mu\nu\lambda}^{(300)} D x^{\mu} D x^{\nu} D x^{\lambda} \\
& \left. + \dots \right],
\end{aligned}$$

$$D = \frac{1}{e} \frac{d}{d\tau} = \text{covar. derivative}$$

nonlocal particle  $\simeq$

open string with spacetime metric  $g_{\mu\nu} \rightarrow 0$

If analogous to String Field Theory,

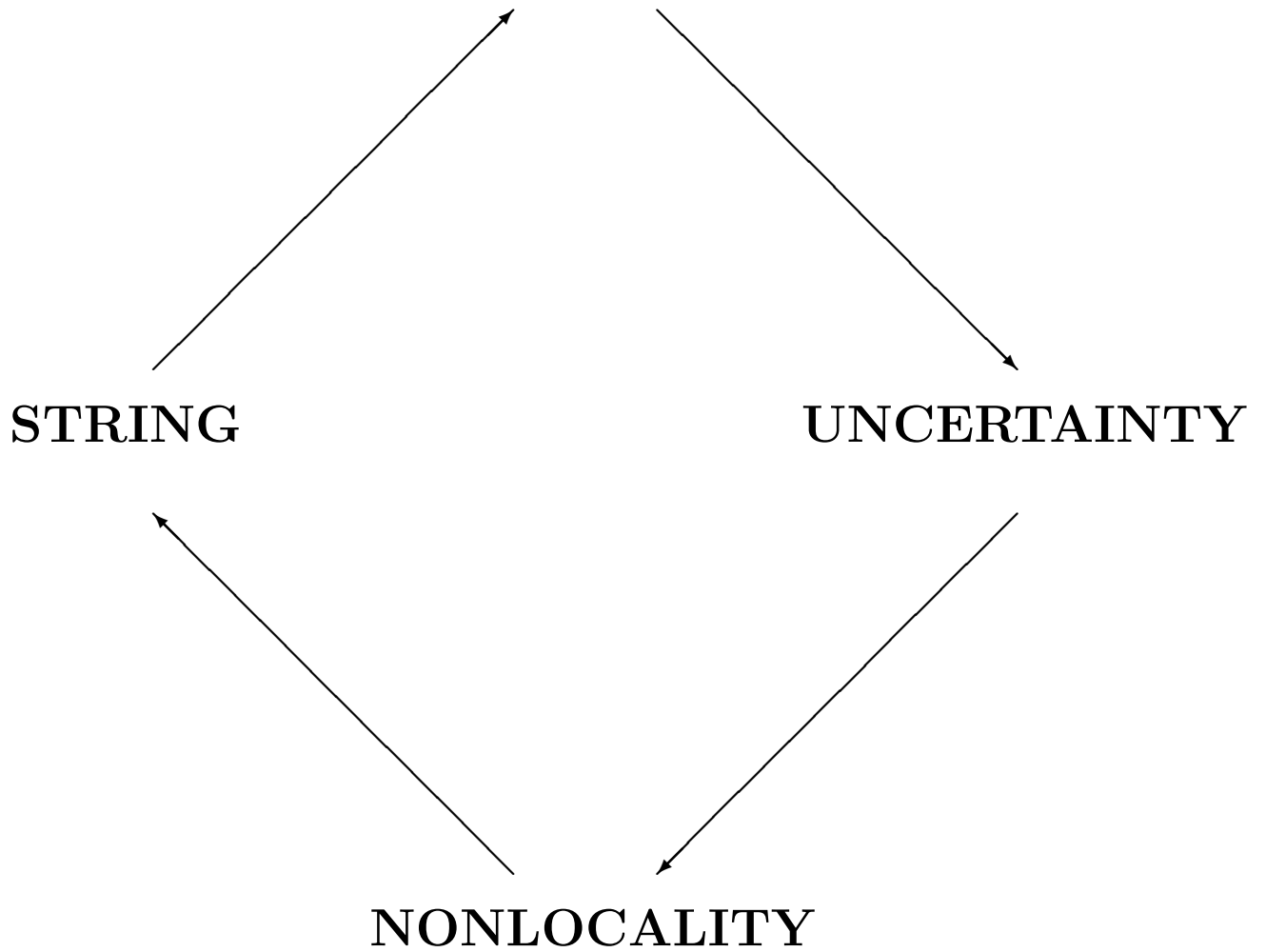
Conformal Symmetry  $\rightarrow$

theory of the background fields  $A^{(\cdot)}(x)$

String Theory  $\simeq$

theory of all theories of particles

**NONCOMMUTATIVITY**



***NONASSOCIATIVITY?***